

ON THE NONLINEAR ULTRA-HYPERBOLIC  
HEAT EQUATION RELATED TO THE SPECTRUM

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**Abstract:** In this paper, we study the nonlinear equation of the form

$$\frac{\partial}{\partial t} u(x, t) - c^2 \square u(x, t) = f(x, t, u(x, t)),$$

where  $\square$  is the *ultra-hyperbolic operator* defined by

$$\square = \left( \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} + \cdots + \frac{\partial^2}{\partial x_p^2} - \frac{\partial^2}{\partial x_{p+1}^2} - \frac{\partial^2}{\partial x_{p+2}^2} - \cdots - \frac{\partial^2}{\partial x_{p+q}^2} \right),$$

$p + q = n$  is the dimension of the Euclidean space  $\mathbb{R}^n$ ,  $(x, t) = (x_1, x_2, \dots, x_n, t) \in \mathbb{R}^n \times (0, \infty)$  and  $c$  is a positive constant.

On the suitable conditions for  $f$ ,  $u$  and for the spectrum of the heat kernel, we can find the unique solution in the compact subset of  $\mathbb{R}^n \times (0, \infty)$ . Moreover, if we put  $q = 0$  we obtain the solution of nonlinear equation related to the heat equation.

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**Key Words:** nonlinear equation, ultra-hyperbolic heat equation

### 1. Introduction

It is well known that for the heat equation

$$\frac{\partial}{\partial t} u(x, t) = c^2 \Delta u(x, t) \quad (1.1)$$

with the initial condition

$$u(x, 0) = f(x),$$

where  $\Delta = \sum_{i=1}^n \frac{\partial^2}{\partial x_i^2}$  is the Laplace operator and  $(x, t) = (x_1, x_2, \dots, x_n, t) \in \mathbb{R}^n \times (0, \infty)$ , and  $f$  is a continuous function, we obtain the solution

$$u(x, t) = \frac{1}{(4c^2\pi t)^{n/2}} \int_{\mathbb{R}^n} \exp\left[-\frac{|x-y|^2}{4c^2t}\right] f(y) dy \quad (1.2)$$

as the solution of (1.1).

Now, (1.2) can be written as  $u(x, t) = E(x, t) * f(x)$ , where

$$E(x, t) = \frac{1}{(4c^2\pi t)^{n/2}} \exp\left[-\frac{|x|^2}{4c^2t}\right]. \quad (1.3)$$

$E(x, t)$  is called *the heat kernel*, where  $|x|^2 = x_1^2 + x_2^2 + \dots + x_n^2$  and  $t > 0$ , see [1, pp. 208-209]. Moreover, we obtain  $E(x, t) \rightarrow \delta$  as  $t \rightarrow 0$ , where  $\delta$  is the Dirac-delta distribution.

We also have studied the generalization of (1.1) in the form

$$\frac{\partial}{\partial t} u(x, t) = -c^2 (-\Delta)^k u(x, t) \quad (1.4)$$

with the initial condition

$$u(x, 0) = f(x),$$

where the operator  $\Delta^k$  denotes *the Laplace operator iterated k-times*. This operator defined as follows

$$\Delta^k = \left( \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} + \dots + \frac{\partial^2}{\partial x_n^2} \right)^k.$$

We obtain  $u(x, t) = E(x, t) * f(x)$  as a solution of (1.4), where (see [4])

$$E(x, t) = \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} \exp\left[-c^2 t \left( \sum_{i=1}^n \xi_i^2 \right)^k + i(\xi, x)\right] d\xi. \quad (1.5)$$

Moreover, we obtain  $E(x, t) \rightarrow \delta$  as  $t \rightarrow 0$ , where  $\delta$  is the Dirac-delta distribution.

Furthermore, we also studied the equation

$$\frac{\partial}{\partial t} u(x, t) = c^2 \square u(x, t), \tag{1.6}$$

where  $\square$  is the ultra-hyperbolic operator, defined by

$$\square = \left( \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} + \dots + \frac{\partial^2}{\partial x_p^2} - \frac{\partial^2}{\partial x_{p+1}^2} - \frac{\partial^2}{\partial x_{p+2}^2} - \dots - \frac{\partial^2}{\partial x_{p+q}^2} \right).$$

We obtain the ultra-hyperbolic heat kernel

$$E(x, t) = \frac{(i)^q}{(4c^2\pi t)^{n/2}} \exp \left[ \frac{\sum_{i=1}^p x_i^2 - \sum_{j=p+1}^{p+q} x_j^2}{4c^2 t} \right],$$

where  $p + q = n$  is the dimension of the Euclidean space  $\mathbb{R}^n$  and  $i = \sqrt{-1}$ . For finding the kernel  $E(x, t)$ , see [5].

Alternatively, we also studied  $E(x, t)$  in the spectrum  $\Omega \subset \mathbb{R}^n$  then we obtain (see [6])

$$E(x, t) = \frac{1}{(2\pi)^n} \int_{\Omega} \exp \left[ c^2 t \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) + i(\xi, x) \right] d\xi.$$

Moreover, we obtain  $E(x, t) \rightarrow \delta$  as  $t \rightarrow 0$ , where  $\delta$  is the Dirac-delta distribution.

In this paper, we extend (1.6) to be the general of the nonlinear form

$$\frac{\partial}{\partial t} u(x, t) - c^2 \square u(x, t) = f(x, t, u(x, t)) \tag{1.7}$$

for  $(x, t) \in \mathbb{R}^n \times (0, \infty)$  and with the following conditions on  $u$  and  $f$  as follows:

- (1)  $u(x, t) \in \mathcal{C}''(\mathbb{R}^n)$  for any  $t > 0$  where  $\mathcal{C}''(\mathbb{R}^n)$  is the space of continuous function with second derivatives.
- (2)  $f$  satisfies the Lipchitz condition, that is

$$|f(x, t, u) - f(x, t, w)| \leq A|u - w|, \quad \text{where } A \text{ is constant and } 0 < A < 1.$$

- (3)

$$\int_0^\infty \int_{\mathbb{R}^n} |f(x, t, u(x, t))| dx dt < \infty$$

for  $x = (x_1, x_2, \dots, x_n) \in \mathbb{R}^n$ ,  $t \in (0, \infty)$  and  $u(x, t)$  is continuous function on  $\mathbb{R}^n \times (0, \infty)$ .

Under such conditions of  $f$ ,  $u$  and for spectrum of  $E(x, t)$ , we obtain the convolution

$$u(x, t) = E(x, t) * f(x, t, u(x, t))$$

as a unique solution in the compact subset of  $\mathbb{R}^n \times (0, \infty)$  and  $E(x, t)$  is an elementary solution defined by (2.5).

## 2. Preliminaries

**Definition 2.1.** Let  $f(x) \in \mathbb{L}_1(\mathbb{R}^n)$ -the space of integrable function in  $\mathbb{R}^n$ . The Fourier transform of  $f(x)$  is defined by

$$\widehat{f}(\xi) = \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{-i(\xi, x)} f(x) dx, \quad (2.1)$$

where  $\xi = (\xi_1, \xi_2, \dots, \xi_n)$ ,  $x = (x_1, x_2, \dots, x_n) \in \mathbb{R}^n$ ,  $(\xi, x) = \xi_1 x_1 + \xi_2 x_2 + \dots + \xi_n x_n$  is the usual inner product in  $\mathbb{R}^n$  and  $dx = dx_1 dx_2 \dots dx_n$ .

Also, the inverse of Fourier transform is defined by

$$f(x) = \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{i(\xi, x)} \widehat{f}(\xi) d\xi. \quad (2.2)$$

**Definition 2.2.** The spectrum of the kernel  $E(x, t)$  defined by (2.5) is the bounded support of the Fourier transform  $\widehat{E}(\xi, t)$  for any fixed  $t > 0$ .

**Definition 2.3.** Let  $\xi = (\xi_1, \xi_2, \dots, \xi_n)$  be a point in  $\mathbb{R}^n$  and we write

$$u = \xi_1^2 + \xi_2^2 + \dots + \xi_p^2 - \xi_{p+1}^2 - \xi_{p+2}^2 - \dots - \xi_{p+q}^2, \quad p + q = n.$$

Denote by  $\Gamma_+ = \{\xi \in \mathbb{R}^n : \xi_1 > 0 \text{ and } u > 0\}$  the set of an interior of the forward cone, and  $\overline{\Gamma}_+$  denotes the closure of  $\Gamma_+$ .

Let  $\Omega$  be spectrum of  $E(x, t)$  defined by Definition 2.2 for any fixed  $t > 0$  and  $\Omega \subset \overline{\Gamma}_+$ . Let  $\widehat{E}(\xi, t)$  be the Fourier transform of  $E(x, t)$  and define

$$\widehat{E}(\xi, t) = \begin{cases} \frac{1}{(2\pi)^{n/2}} \exp \left[ c^2 t \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) \right] & \text{for } \xi \in \Gamma_+, \\ 0 & \text{for } \xi \notin \Gamma_+. \end{cases} \quad (2.3)$$

**Lemma 2.1.** Let  $L$  be the operator defined by

$$L = \frac{\partial}{\partial t} - c^2 \square, \quad (2.4)$$

where  $\square$  is the ultra-hyperbolic operator defined by

$$\square = \left( \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} + \cdots + \frac{\partial^2}{\partial x_p^2} - \frac{\partial^2}{\partial x_{p+1}^2} - \frac{\partial^2}{\partial x_{p+2}^2} - \cdots - \frac{\partial^2}{\partial x_{p+q}^2} \right),$$

$p + q = n$  is the dimension of  $\mathbb{R}^n$ ,  $(x_1, x_2, \dots, x_n) \in \mathbb{R}^n$ ,  $t \in (0, \infty)$  and  $c$  is a positive constant. Then we obtain

$$E(x, t) = \frac{1}{(2\pi)^n} \int_{\Omega} \exp \left[ c^2 t \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) + i(\xi, x) \right] d\xi \quad (2.5)$$

as a elementary solution of (2.4) in the spectrum  $\Omega \subset \mathbb{R}^n$  for  $t > 0$ .

*Proof.* Let  $LE(x, t) = \delta(x, t)$ , where  $E(x, t)$  is the kernel or the elementary solution of operator  $L$  and  $\delta$  is the Dirac-delta distribution. Thus

$$\frac{\partial}{\partial t} E(x, t) - c^2 \square E(x, t) = \delta(x) \delta(t).$$

Take the Fourier transform defined by (2.1) to both sides of the equation, we obtain

$$\frac{\partial}{\partial t} \widehat{E(\xi, t)} - c^2 \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) \widehat{E(\xi, t)} = \frac{1}{(2\pi)^{n/2}} \delta(t).$$

Thus

$$\widehat{E(\xi, t)} = \frac{H(t)}{(2\pi)^{n/2}} \exp \left[ c^2 t \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) \right],$$

where  $H(t)$  is the Heaviside function. Since  $H(t) = 1$  for  $t > 0$ . Therefore,

$$\widehat{E(\xi, t)} = \frac{1}{(2\pi)^{n/2}} \exp \left[ c^2 t \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) \right]$$

which has been already defined by (2.3). Thus

$$E(x, t) = \frac{1}{(2\pi)^{n/2}} \int_{\mathbb{R}^n} e^{i(\xi, x)} \widehat{E(\xi, t)} d\xi = \frac{1}{(2\pi)^{n/2}} \int_{\Omega} e^{i(\xi, x)} \widehat{E(\xi, t)} d\xi,$$

where  $\Omega$  is the spectrum of  $E(x, t)$ . Thus from (2.3)

$$E(x, t) = \frac{1}{(2\pi)^n} \times \int_{\Omega} \exp \left[ c^2 t \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) + i(\xi, x) \right] d\xi \quad \text{for } t > 0. \quad \square$$

**Definition 2.4.** Let us extend  $E(x, t)$  to  $\mathbb{R}^n \times \mathbb{R}$  by setting

$$E(x, t) = \begin{cases} \frac{1}{(2\pi)^n} \int_{\Omega} \exp \left[ c^2 t \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) + i(\xi, x) \right] d\xi & \text{for } t > 0, \\ 0 & \text{for } t \leq 0, \end{cases}$$

### 3. Main Results

**Theorem 3.1.** *The kernel  $E(x, t)$  defined by (2.5) have the following properties:*

(1)  $E(x, t) \in \mathcal{C}^\infty$ -the space of continuous function for  $x \in \mathbb{R}^n$ ,  $t > 0$  with infinitely differentiable.

(2)  $\left( \frac{\partial}{\partial t} - c^2 \square \right) E(x, t) = 0$  for  $t > 0$ .

(3)  $|E(x, t)| \leq \frac{2^{2-n}}{\pi^{n/2}} \frac{M(t)}{\Gamma(\frac{p}{2})\Gamma(\frac{q}{2})}$ , for  $t > 0$ , where  $M(t)$  is a function of  $t$  in the spectrum  $\Omega$  and  $\Gamma$  denote the Gamma function. Thus  $E(x, t)$  is bounded for any fixed  $t > 0$ .

(4)  $\lim_{t \rightarrow 0} E(x, t) = \delta$ .

*Proof.* (1) From (2.5), since

$$\frac{\partial^n}{\partial x^n} E(x, t) = \frac{1}{(2\pi)^n} \int_{\Omega} \frac{\partial^n}{\partial x^n} \exp \left[ c^2 t \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) + i(\xi, x) \right] d\xi.$$

Thus  $E(x, t) \in \mathcal{C}^\infty$  for  $x \in \mathbb{R}^n$ ,  $t > 0$ .

(2) By computing directly, we obtain

$$\left( \frac{\partial}{\partial t} - c^2 \square \right) E(x, t) = 0.$$

(3) We have

$$E(x, t) = \frac{1}{(2\pi)^n} \int_{\Omega} \exp \left[ c^2 t \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) + i(\xi, x) \right] d\xi,$$

$$|E(x, t)| \leq \frac{1}{(2\pi)^n} \int_{\Omega} \exp \left[ c^2 t \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) \right] d\xi.$$

By changing to bipolar coordinates

$$\xi_1 = r\omega_1, \xi_2 = r\omega_2, \dots, \xi_p = r\omega_p \quad \text{and}$$

$$\xi_{p+1} = s\omega_{p+1}, \xi_{p+2} = s\omega_{p+2}, \dots, \xi_{p+q} = s\omega_{p+q},$$

where  $\sum_{i=1}^p \omega_i^2 = 1$  and  $\sum_{j=p+1}^{p+q} \omega_j^2 = 1$ . Thus

$$|E(x, t)| \leq \frac{1}{(2\pi)^n} \int_{\Omega} \exp [c^2 t (s^2 - r^2)] r^{p-1} s^{q-1} dr ds d\Omega_p d\Omega_q,$$

where  $d\xi = r^{p-1} s^{q-1} dr ds d\Omega_p d\Omega_q$ ,  $d\Omega_p$  and  $\Omega_q$  are the elements of surface area of the unit sphere in  $\mathbb{R}^p$  and  $\mathbb{R}^q$  respectively. Since  $\Omega \subset \mathbb{R}^n$  is the spectrum of  $E(x, t)$  and we suppose  $0 \leq r \leq R$  and  $0 \leq s \leq L$ , where  $R$  and  $L$  are constants. Thus we obtain

$$\begin{aligned} |E(x, t)| &\leq \frac{\Omega_p \Omega_q}{(2\pi)^n} \int_0^R \int_0^L \exp [c^2 t (s^2 - r^2)] r^{p-1} s^{q-1} ds dr \\ &= \frac{\Omega_p \Omega_q}{(2\pi)^n} M(t) \quad \text{for any fixed } t > 0 \text{ in the spectrum } \Omega \\ &= \frac{2^{2-n}}{\pi^{n/2}} \frac{M(t)}{\Gamma(\frac{p}{2})\Gamma(\frac{q}{2})}, \end{aligned} \quad (3.1)$$

where

$$M(t) = \int_0^R \int_0^L \exp [c^2 t (s^2 - r^2)] r^{p-1} s^{q-1} ds dr \quad (3.2)$$

is a function of  $t > 0$ ,  $\Omega_p = \frac{2\pi^{p/2}}{\Gamma(\frac{p}{2})}$  and  $\Omega_q = \frac{2\pi^{q/2}}{\Gamma(\frac{q}{2})}$ . Thus, for any fixed  $t > 0$ ,  $E(x, t)$  is bounded.

(4) By (2.5), we have

$$E(x, t) = \frac{1}{(2\pi)^n} \int_{\Omega} \exp \left[ c^2 t \left( \sum_{j=p+1}^{p+q} \xi_j^2 - \sum_{i=1}^p \xi_i^2 \right) + i(\xi, x) \right] d\xi.$$

Since  $E(x, t)$  exists, then

$$\lim_{t \rightarrow 0} E(x, t) = \frac{1}{(2\pi)^n} \int_{\Omega} e^{i(\xi, x)} d\xi = \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i(\xi, x)} d\xi = \delta(x), \quad \text{for } x \in \mathbb{R}^n.$$

See [2, p. 396, equation (10.2.19b)].  $\square$

**Theorem 3.2.** *Given the nonlinear equation*

$$\frac{\partial}{\partial t} u(x, t) - c^2 \square u(x, t) = f(x, t, u(x, t)) \quad (3.3)$$

for  $(x, t) \in \mathbb{R}^n \times (0, \infty)$  and with the following conditions on  $u$  and  $f$  as follows:

(1)  $u(x, t) \in \mathcal{C}''(\mathbb{R}^n)$  for any  $t > 0$ , where  $\mathcal{C}''(\mathbb{R}^n)$  is the space of continuous function with second derivatives.

(2)  $f$  satisfies the Lipchitz condition, that is

$$|f(x, t, u) - f(x, t, w)| \leq A|u - w|, \quad \text{where } A \text{ is constant and } 0 < A < 1.$$

(3)  $\int_0^\infty \int_{\mathbb{R}^n} |f(x, t, u(x, t))| dx dt < \infty$ , for  $x = (x_1, x_2, \dots, x_n) \in \mathbb{R}^n$ ,  $t \in (0, \infty)$  and  $u(x, t)$  is continuous function on  $\mathbb{R}^n \times (0, \infty)$ .

Then, for the spectrum of  $E(x, t)$  we obtain the convolution

$$u(x, t) = E(x, t) * f(x, t, u(x, t)) \quad (3.4)$$

as a unique solution of (3.3) for  $x \in \Omega_0$ , where  $\Omega_0$  is an compact subset of  $\mathbb{R}^n$  and  $0 \leq t \leq T$  with  $T$  is constant and  $E(x, t)$  is an elementary solution defined by (2.5) and also  $u(x, t)$  is bounded.

In particular, if we put  $q = 0$  in (3.3) then (3.3) reduces to the nonlinear heat equation.

*Proof.* Convolving both sides of (3.3) with  $E(x, t)$  and then we obtain the solution

$$u(x, t) = E(x, t) * f(x, t, u(x, t)),$$

or

$$u(x, t) = \int_{-\infty}^{\infty} \int_{\mathbb{R}^n} E(r, s) f(x - r, t - s, u(x - r, t - s)) dr ds,$$

where  $E(r, s)$  is given by Definition 2.4.

We next show that  $u(x, t)$  is bounded on  $\mathbb{R}^n \times (0, \infty)$ . We have

$$\begin{aligned} |u(x, t)| &\leq \int_{-\infty}^{\infty} \int_{\mathbb{R}^n} |E(r, s)| |f(x - r, t - s, u(x - r, t - s))| dr ds \\ &\leq \frac{2^{2-n} N.M(t)}{\pi^{n/2} \Gamma(\frac{p}{2}) \Gamma(\frac{q}{2})} \end{aligned}$$

by the condition (3) and (3.1), where

$$N = \int_0^\infty \int_{\mathbb{R}^n} |f(x, t, u(x, t))| dx dt.$$



Thus  $u(x, t)$  is bounded on  $\mathbb{R}^n \times (0, \infty)$ .

To show that  $u(x, t)$  is unique, suppose there is another solution  $w(x, t)$  of equation (3.3). Let the operator

$$L = \frac{\partial}{\partial t} - c^2 \square$$

then (3.3) can be written in the form

$$L u(x, t) = f(x, t, u(x, t)).$$

Thus

$$L u(x, t) - L w(x, t) = f(x, t, u(x, t)) - f(x, t, w(x, t)).$$

By the condition (2) of the theorem,

$$|L u(x, t) - L w(x, t)| \leq A |u(x, t) - w(x, t)|. \tag{3.5}$$

Let  $\Omega_0 \times (0, T]$  be compact subset of  $\mathbb{R}^n \times (0, \infty)$  and  $L : \mathcal{C}''(\Omega_0) \rightarrow \mathcal{C}''(\Omega_0)$  for  $0 \leq t \leq T$ .

Now  $(\mathcal{C}''(\Omega_0), \|\cdot\|)$  is a Banach space, where  $u(x, t) \in \mathcal{C}''(\Omega_0)$  for  $0 \leq t \leq T$ ,  $\|\cdot\|$  given by

$$\|u(x, t)\| = \sup_{x \in \Omega_0} |u(x, t)|.$$

Then, from (3.5) with  $0 < A < 1$ , the operator  $L$  is a contraction mapping on  $\mathcal{C}''(\Omega_0)$ . Since  $(\mathcal{C}''(\Omega_0), \|\cdot\|)$  is a Banach space and  $L : \mathcal{C}''(\Omega_0) \rightarrow \mathcal{C}''(\Omega_0)$  is a contraction mapping on  $\mathcal{C}''(\Omega_0)$ , by Contraction Theorem, see [3, p. 300], we obtain the operator  $L$  has a fixed point and has uniqueness property.

Thus  $u(x, t) = w(x, t)$ . It follows that the solution  $u(x, t)$  of (3.3) is unique for  $(x, t) \in \Omega_0 \times (0, T]$ , where  $u(x, t)$  is defined by (3.4).

In particular, if we put  $q = 0$  in (3.3) then (3.3) reduces to the nonlinear heat equation

$$\frac{\partial}{\partial t} u(x, t) - c^2 \Delta u(x, t) = f(x, t, u(x, t))$$

which has solution

$$u(x, t) = E(x, t) * f(x, t, u(x, t)),$$

where  $E(x, t)$  is defined by (2.5) with  $q = 0$ . That is complete of proof. □

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### References

- [1] F. John, *Partial Differential Equations*, Applied Math. Sc., Volume 4, Springer Verlag (1982).
- [2] R. Haberman, *Elementary Applied Partial Differential Equations*, Second Edition, Prentice Hall International Inc. (1983).
- [3] E. Kreyszig, *Introductory Functional Analysis with Applications*, John Wiley and Sons Inc. (1978).
- [4] K. Nonlaopon, A. Kananthai, On the generalized heat kernel, *Computational Technologies*, **9**, No. 1 (2004), 3-10.
- [5] K. Nonlaopon, A. Kananthai, On the ultrahyperbolic heat kernel, *International Journal of Applied Mathematics*, **13**, No. 2 (2003), 215-225.
- [6] K. Nonlaopon, A. Kananthai, On the ultra-hyperbolic heat kernel related to the spectrum, *International Journal of Pure and Applied Mathematics*, **17**, No. 1, 19-28.