

ASYMPTOTIC EXPRESSIONS FOR THE EIGENVALUES AND
EIGENVECTORS OF A SYSTEM OF SECOND ORDER
DIFFERENTIAL EQUATIONS WITH
A TURNING POINT (EXTENSION-1)

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Abstract: Consider the system of second order differential equation

$$y''(x) + (\lambda^2 R(x) + Q(x))y(x) = 0, \quad 0 \leq x \leq \pi,$$

where $y(x) = (y_1(x), y_2(x))^T$, $Q(x) = \begin{pmatrix} p(x) & r(x) \\ r(x) & q(x) \end{pmatrix}$, $R(x) = \begin{pmatrix} s(x) & 0 \\ 0 & t(x) \end{pmatrix}$, $s(x) = x^m s_1(x)$, $t(x) = x^m t_1(x)$, $s_1(x) > 0$, $t_1(x) > 0$ and $p(x)$, $q(x)$, $r(x)$, $s_1(x)$, $t_1(x)$ are real-valued functions having continuous second order derivatives at x , $0 \leq x \leq \pi$, m being a real constant and λ , a real parameter.

In the present paper we consider m , a negative real number and determine the asymptotic solutions alongwith their derivatives for such a system for large values of the parameter λ and apply these to determine the asymptotic expressions for the distribution of the eigenvalues and the normalized eigenvectors under the Dirichlet boundary conditions.

AMS Subject Classification: 35B40, 37K40

Key Words: asymptotic solutions, turning points, Dirichlet boundary conditions, normalized eigenvectors

1. Introduction

In the paper of Sengupta [4], the system of second order differential equations

considered was

$$y''(x) + (\lambda^2 R(x) + Q(x))y(x) = 0, \quad 0 \leq x \leq \pi, \quad (1)$$

where $y(x) = (y_1(x), y_2(x))^T$, $Q(x) = \begin{pmatrix} p(x) & r(x) \\ r(x) & q(x) \end{pmatrix}$, $R(x) = \begin{pmatrix} s(x) & 0 \\ 0 & t(x) \end{pmatrix}$, $p(x)$, $q(x)$, $r(x)$, are real-valued functions having continuous second order derivatives at x , $0 \leq x \leq \pi$, and $s(x) = x^m s_1(x)$, $t(x) = x^m t_1(x)$, $s_1(x) > 0$, $t_1(x) > 0$ have continuous second order derivatives at x , $0 \leq x \leq \pi$, m being positive constant and λ , a real parameter.

In the present paper for the system (1) we consider m as a negative constant and in particular we choose $m = -1$ i.e., we consider the system (1) where

$$s(x) = \frac{s_1(x)}{x}, \quad t(x) = \frac{t_1(x)}{x}, \quad 0 \leq x \leq \pi, \quad (2)$$

where $s_1(x)$, $t_1(x) > \rho^2 > 0$ are continuous functions of x , $0 \leq x \leq \pi$.

The boundary conditions considered in the problem are the Dirichlet boundary conditions

$$y_1(0) = y_2(0) = y_1(\pi) = y_2(\pi) = 0 \quad (3)$$

satisfied by the solution $y(x) = (y_1(x), y_2(x))^T$ of the system (1)-(2) at $x = 0, x = \pi$. The problem (1)-(3) may be termed after Langer [2], the singular second order system with a simple turning point. Following Neumark [3], we choose $R(x)$ in (1) as a diagonal matrix. As stated explicitly in Dorodnicyn [1], the comparing equations considered for the problem (1)-(2) are the system of Bessel equations

$$\begin{aligned} v_1''(z) + \frac{v_1(z)}{z} &= 0, \\ v_2''(z) + \frac{v_2(z)}{z} &= 0, \end{aligned} \quad (4)$$

whose solutions are

$$\begin{aligned} v_1(z) &= -z^{\frac{1}{2}} Y_1(2z^{\frac{1}{2}}), \\ v_2(z) &= z^{\frac{1}{2}} J_1(2z^{\frac{1}{2}}), \end{aligned} \quad (5)$$

$J_1(\cdot)$, $Y_1(\cdot)$ being the Bessel functions of order one of the first and second kind respectively, have the Wronskian $W(J_1(z), Y_1(z)) = \frac{2}{z}$ (see Watson [5]).

Hence,

$$W(v_1, v_2) = 1, \tag{6}$$

$$\begin{aligned} \text{At } z = 0, v_1(0) &= 0 \text{ (in the sense of limit } v_1(z) \text{ as } z \rightarrow 0) \\ \text{and } v_2(0) &= 0. \end{aligned} \tag{7}$$

It may be noted that $v_1'(0)$ does not exist; in fact $v_1'(z) \rightarrow \infty$, as $z \rightarrow 0$.

In the present paper, for the problem (1)-(3), we determine the asymptotic expressions of the solutions, their derivatives and apply these to determine the asymptotic expressions for the n -th eigenvalue λ_n and the corresponding eigenvector $\Psi(x, \lambda_n)$, for large values of the parameter λ .

As the present work is an extension of the work of Sengupta [4], we mention explicitly only those which are considerably different.

2. Preliminaries

Let us substitute

$$\begin{aligned} w_1(x) &= \left(\frac{1}{2} \int_0^x \sqrt{s(z)} dz\right)^2, \\ w_2(x) &= \left(\frac{1}{2} \int_0^x \sqrt{t(z)} dz\right)^2. \end{aligned} \tag{8}$$

Then

$$\begin{aligned} \sqrt{s(x)} &= \frac{w_1'(x)}{\sqrt{w_1(x)}}, \\ \sqrt{t(x)} &= \frac{w_2'(x)}{\sqrt{w_2(x)}}, \end{aligned} \tag{9}$$

and $w_i'(0) \neq 0, i = 1, 2$, since $x = 0$ is a zero of order one of $w_i(x), i = 1, 2$.

The auxiliary equation is

$$y''(x) + (\lambda^2 R(x) + Q_0(x))y(x) = 0 \tag{10}$$

where $R(x)$ is given in (1) and

$$Q_0(x) = \begin{pmatrix} p_0(x) & 0 \\ 0 & q_0(x) \end{pmatrix},$$

$$\begin{aligned}
 p_0(x) &= -\sqrt{w_1'(x)}D^2(w_1'(x))^{-\frac{1}{2}}. \\
 q_0(x) &= -\sqrt{w_2'(x)}D^2(w_2'(x))^{-\frac{1}{2}}, D \equiv \frac{d}{dx}
 \end{aligned}
 \tag{11}$$

we assume further that

$$|p(x) - p_0(x)|, |q(x) - q_0(x)|, |r(x)| \leq c.x^k \tag{12}$$

where $k > 0$ and c is a non-zero constant. Then adopting the technique used to prove Theorem-1 of Sengupta[4], the following theorem which establishes the existence of a particular solution of the system (10) is proved.

Theorem 1. *Let*

$$y_i(x) = C_{i1}(w_i'(x))^{-\frac{1}{2}}v_1(\lambda^2w_i(x)) + C_{i2}(w_i'(x))^{-\frac{1}{2}}v_2(\lambda^2w_i(x)), \quad i = 1, 2, \tag{13}$$

where $C_{ij}, i, j = 1, 2$ are constants, $(v_1(x), v_2(x))^T$ is the solution of the system (4), satisfy at $x = 0$ the conditions(7) and $w_i(x), i = 1, 2$ are those given in (8).

Then $y(x) = (y_1(x), y_2(x))^T$ satisfies the system (10).

3. Asymptotic Estimates of Certain Integrals

Let $\alpha_i(x) = (\alpha_{i1}(x), \alpha_{i2}(x))^T, i = 1, 2$, where

$$\alpha_{ij}(x) = (w_j'(x))^{-\frac{1}{2}}.v_i(\lambda^2w_j(x)), \quad i, j = 1, 2 \tag{14}$$

Then for $j = 1, 2; \alpha_{1j}(0) = (w_j'(0))^{-\frac{1}{2}}$ (in the sense of limit $\alpha_{1j}(x)$ as $x \rightarrow 0$), $\alpha_{2j}(0) = (w_j'(0))^{-\frac{1}{2}}, \alpha'_{2j}(0) = (w_j'(0))^{-\frac{1}{2}}$ but $\alpha'_{1j}(0)$ does not exist.

As $(v_1(x), v_2(x))^T$ are the solutions of (4), following Theorem 1, it follows that $\alpha_1(x), \alpha_2(x)$ are the two linearly independent solutions of the system (10).

However, the bilinear concomitant $[\alpha_1(x), \alpha_2(x)]$ is independent of x and is a function of λ alone and

$$[\alpha_1(x), \alpha_2(x)] = 2\lambda^2 = \frac{1}{\delta} \quad (\text{say}) \tag{15}$$

Let

$$f(x) = Q(x) - Q_0(x) = (f_{ij}(x)), i, j = 1, 2 \tag{16}$$

where $f_{ij}(x), i, j = 1, 2$ satisfy the conditions (12).

Using the conditions (12), we settle the convergence of the integrals

$$\int_0^x \frac{f_{11}(z)}{\sqrt{s(z)}} dz, \quad \int_0^x \frac{f_{12}(z)}{\sqrt[4]{s(z)t(z)}} dz$$

and

$$\int_0^x \frac{f_{22}(z)}{\sqrt{t(z)}} dz.$$

Following Dorodnicyn [1] and Watson [5] we obtain the following asymptotic relations for large λ :

$$\begin{aligned} \int_0^x \alpha_{11}^2(z) f_{11}(z) dz, \int_0^x \alpha_{21}^2(z) f_{11}(z) dz, \int_0^x \alpha_{11}(z) \alpha_{21}(z) f_{11}(z) dz \\ = c\lambda \cdot \int_0^x \frac{f_{11}(z)}{\sqrt{s(z)}} dz + 0(\lambda^{-1}), \end{aligned}$$

$$\begin{aligned} \int_0^x \alpha_{12}^2(z) f_{22}(z) dz, \int_0^x \alpha_{22}^2(z) f_{22}(z) dz, \int_0^x \alpha_{12}(z) \alpha_{22}(z) f_{22}(z) dz \\ = c\lambda \cdot \int_0^x \frac{f_{22}(z)}{\sqrt{s(z)}} dz + 0(\lambda^{-1}), \end{aligned}$$

$$\begin{aligned} \int_0^x \alpha_{12} \alpha_{21}(z) f_{12}(z) dz, \int_0^x \alpha_{22} \alpha_{11}(z) f_{12}(z) dz, \int_0^x \alpha_{11} \alpha_{12}(z) f_{12}(z) dz, \\ \int_0^x \alpha_{21} \alpha_{22}(z) f_{12}(z) dz = c\lambda \cdot \int_0^x \frac{f_{12}(z)}{\sqrt[4]{s(z)t(z)}} dz + 0(z^{-1}), \quad (17) \end{aligned}$$

where c represents different numerical constants.

4. Asymptotic Relations for the Solutions

Let

$$\alpha(z) = \begin{pmatrix} \alpha_{11}(z) & \alpha_{21}(z) \\ \alpha_{12}(z) & \alpha_{22}(z) \end{pmatrix}, \quad \alpha_i(z) = (\alpha_{i1}(z), \alpha_{i2}(z))^T, \quad i = 1, 2... \quad (18)$$

and

$$\begin{aligned}
 y(x, z) &= \begin{pmatrix} y_{11}(x, z) & y_{21}(x, z) \\ y_{12}(x, z) & y_{22}(x, z) \end{pmatrix} \\
 &= \begin{pmatrix} \alpha_{21}(z)\alpha_{11}(x) - \alpha_{11}(z)\alpha_{21}(x) & \alpha_{22}(z)\alpha_{11}(x) - \alpha_{12}(z)\alpha_{21}(x) \\ \alpha_{21}(z)\alpha_{12}(x) - \alpha_{11}(z)\alpha_{22}(x) & \alpha_{22}(z)\alpha_{12}(x) - \alpha_{12}(z)\alpha_{22}(x) \end{pmatrix}, \quad (19)
 \end{aligned}$$

where $\alpha_{ij}(\cdot)$ are given by (14).

Further let

$$\beta(z) = \begin{pmatrix} \beta_{11}(z) & \beta_{21}(z) \\ \beta_{12}(z) & \beta_{22}(z) \end{pmatrix}, \quad \beta_i(z) = (\beta_{i1}(z), \beta_{i2}(z))^T, \quad i = 1, 2, \dots \quad (20)$$

Then following Sengupta [4], the following theorem is proved.

Theorem 2. *Let the elements of the matrix $f(x)$, defined by (16), satisfy the conditions (12). Then the two linearly independent solutions of (1) are given by*

$$y_i(x) = (y_{i1}(x), y_{i2}(x))^T = \alpha_i(x) + \beta_i(x), \quad i = 1, 2, \dots, \quad (21)$$

where $\alpha_i(x), i = 1, 2$, given by (14), are the linearly independent solutions of (10) and $\beta_i; i = 1, 2$ defined by (20) is the unique solution of the integral equation

$$\beta(x) = \delta \int_0^x y(x, z) f(z) (\alpha(z) + \beta(z)) dz \quad (22)$$

$y(x, z), \delta$ being those defined in (19), (15) respectively.

By using the asymptotic relations as given in Section 3, we obtain for large λ that

$$\beta(x) = \lambda^{-1} \alpha(x) MP(x) + o(\lambda^{-2} \|\alpha(x)\|), \quad (23)$$

where $M = \begin{pmatrix} \frac{1}{2} & \frac{1}{2\pi} \\ -\frac{\pi}{2} & -\frac{1}{2} \end{pmatrix}$, and

$$P(x) = \int_0^x \left[\frac{f_{11}(z)}{\sqrt{s(z)}} + 2 \frac{f_{12}(z)}{\sqrt[4]{s(z)t(z)}} + \frac{f_{22}(z)}{\sqrt{t(z)}} \right] dz \quad (24)$$

Hence,

$$y(x) = \alpha(x) + \lambda^{-1} \alpha(x) MP(x) + o(\lambda^{-2} \|\alpha(x)\|) \quad (25)$$

Differentiating with respect to x first and then using the asymptotic estimates of the integrals as given in Section 3, it follows from (22) that for large λ ,

$$\beta'(x) \cong \delta y(x, x)f(x)\alpha(x) + \lambda^{-1}\alpha'(x)MP(x) + 0(\lambda^{-2}||\alpha(x)||), \tag{26}$$

where

$$y(x, x) = (\alpha_{11}(x)\alpha_{22}(x) - \alpha_{12}(x)\alpha_{21}(x)).I_1, I_1 = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \tag{27}$$

and $\delta, M, p(x)$ are those given before. Hence,

$$y'(x) = \alpha'(x) + \delta y(x, x)f(x)\alpha(x) + \lambda^{-1}\alpha'(x)MP(x) + 0(\lambda^{-2}.||\alpha'(x)||) \tag{28}$$

5. Asymptotic Representation for the Eigenvalues

The four linearly independent solutions of the system (1)-(2) are

$$\begin{aligned} U_1(x) &= (U_{11}(x), U_{12}(x))^T = (\alpha_{11}(x) + B_{11}(x), B_{12}(x))^T, \\ U_2(x) &= (U_{21}(x), U_{22}(x))^T = (\alpha_{21}(x) + B_{21}(x), B_{22}(x))^T, \\ \bar{U}_3(x) &= (\bar{U}_{11}(x), \bar{U}_{12}(x))^T = (\bar{B}_{11}(x), \alpha_{12}(x) + \bar{B}_{12}(x))^T \\ &\hspace{15em} \text{and} \\ \bar{U}_4(x) &= (\bar{U}_{21}(x), \bar{U}_{22}(x))^T = (\bar{B}_{21}(x), \alpha_{22}(x) + \bar{B}_{22}(x))^T, \end{aligned} \tag{29}$$

where the $U_{ij}(x), \bar{U}_{ij}(x), B_{ij}(x), \bar{B}_{ij}(x)$ involving the $\alpha_{ij}(x)$, $i, j = 1, 2$ are defined in the same way as given in Section-5 of Sengupta [4].

Following Sengupta [4], it readily follows that as $\lambda_n \rightarrow \infty$,

$$\begin{aligned} U_{1j}(0, \lambda_n), \bar{U}_{1j}(0, \lambda_n) &= (w'_j(0))^{-\frac{1}{2}} + 0(\lambda_n^{-\frac{1}{2}}), \\ U_{2j}(0, \lambda_n), \bar{U}_{2j}(0, \lambda_n) &= 0(\lambda_n^{-2}), \end{aligned}$$

$$\begin{aligned} U_{1j}(\pi, \lambda_n), \bar{U}_{1j}(\pi, \lambda_n) \\ = \lambda_n^{\frac{1}{2}}(w_j(\pi))^{\frac{1}{4}}(w'_j(\pi))^{-\frac{1}{2}}(\pi)\pi^{\frac{1}{2}}Cos(2\lambda_n(w_j(\pi))^{\frac{1}{2}} - \frac{\pi}{4}) + 0(\lambda_n^{-\frac{1}{2}}), \end{aligned}$$

$$\begin{aligned} U_{2j}(\pi, \lambda_n), \bar{U}_{2j}(\pi, \lambda_n) \\ = \lambda_n^{\frac{1}{2}}(w_j(\pi))^{\frac{1}{4}}(w'_j(\pi))^{-\frac{1}{2}}(\pi)^{\frac{1}{2}}Sin(2\lambda_n(w_j(\pi))^{\frac{1}{2}} - \frac{\pi}{4}) + 0(\lambda_n^{-\frac{1}{2}}), \end{aligned}$$

where $i, j = 1, 2$. and

$$\beta_{ij}(0), \bar{\beta}_{ij}(0), \beta_{ij}(\pi), \bar{\beta}_{ij}(\pi) = 0(1), i, j = 1, 2 \tag{30}$$

Hence from the general solution

$$V(x) = AU_1(x) + BU_2(x) + CU_3(x) + DU_4(x),$$

A, B, C, D constants, of the system (1) by using the Dirichlet boundary conditions (3) and the relations (30) it follows in the same way as in Sengupta [4] that the asymptotic expressions for the eigenvalues (λ_n) are given by

$$\lambda_n = \frac{2(n + \frac{1}{4})\pi}{\int_0^\pi \sqrt{s(x)}dx} + 0(n^{-2}) \tag{31}$$

or,

$$\lambda_n = \frac{2(n + \frac{1}{4})\pi}{\int_0^\pi \sqrt{t(x)}dx} + 0(n^{-2}) \tag{32}$$

as $n \rightarrow \infty$, according as

$$\int_0^\pi \sqrt{s(x)}dx > or < \int_0^\pi \sqrt{t(x)}dx$$

It is to be noted that the asymptotic expressions of the eigenvalues λ_n in the present problem with Neumann boundary conditions (i.e., $y'_1(0) = y'_2(0) = y'_1(\pi) = y'_2(\pi) = 0$ satisfied by the solution $y(x) = (y_1(x), y_2(x))^T$ of the system (1) at $x = 0, x = \pi$) cannot be derived, as already shown that $v'_1(x) \rightarrow \infty$ as $x \rightarrow 0$. This has also been pointed out by Dorodnicyn [1], (p.32).

6. Asymptotic Representation of the Normalized Eigenvector

Arguing in the same way as in Sengupta [4], it follows that the normalized eigenvector $\Psi(x, \lambda_n)$ with normalizing constant A_n is given by

$$\Psi(x, \lambda_n) = A_n^{-\frac{1}{2}} [a_{2n}U_1(x, \lambda_n) - a_{1n}\bar{U}_3(x, \lambda_n)] \tag{33}$$

where

$$A_n = a_{1n}a_{2n}(a_{1n} + a_{2n}),$$

$$a_{1n} = \int_0^\pi s(x)U_{11}^2(x)dx$$

$$a_{2n} = \int_0^\pi t(x)\overline{U}_{12}^2(x)dx$$

Since $U_1(x), \overline{U}_3(x), j = 1, 2$ behave asymptotically in the same way as $y_j(x), j = 1, 2$ of Theorem 2 it follows that as $\lambda \rightarrow \infty$.

$$a_{1n} = \lambda_n \cdot 2\pi(w_1(\pi))^{\frac{1}{2}} + 0(1),$$

$$a_{2n} = \lambda_n \cdot 2\pi(w_2(\pi))^{\frac{1}{2}} + 0(1) \tag{34}$$

Also, from (14), $\alpha_{2j} = 0(\lambda_n)^{\frac{1}{2}}$, for $v_2(z) \cong z^{\frac{1}{4}} \cdot \pi^{-\frac{1}{2}}$, follows from the asymptotic formula for the Bessel function $J_1(\cdot)$. Hence,

$$A_n^{-\frac{1}{2}} = \lambda_n^{-\frac{3}{2}}(2\pi)^{-\frac{3}{2}}(w_1(\pi)w_2(\pi))^{-\frac{1}{4}}(w_1^{\frac{1}{2}}(\pi) + w_2^{\frac{1}{2}}(\pi))^{-\frac{1}{2}}(1 + 0(\lambda_n^{-1})) \tag{35}$$

Thus, the asymptotic expressions for normalized eigenvector corresponding to the eigenvalue n for the Dirichlet boundary conditions is given by

$$\Psi(x, \lambda_n) = \lambda_n^{-\frac{1}{2}}(2\pi)^{\frac{1}{2}}(w_1(\pi)w_2(\pi))^{-\frac{1}{4}} \cdot (w_1^{\frac{1}{2}}(\pi) + w_2^{\frac{1}{2}}(\pi))$$

$$(w_2^{\frac{1}{2}}(\pi)\alpha_{11}(x), -w_1^{\frac{1}{2}}(\pi)\alpha_{12}(x))(1 + 0(\lambda_n^{-1})), \tag{36}$$

as $\lambda_n \rightarrow \infty$.

Acknowledgments

The author expresses his grateful thanks to Prof. (Dr.) N. K. Chakravarty, U. G. C. Professor (Retd.) Dept. of Pure Math. Calcutta University for the extremely helpful and constructive suggestions during the preparation of the present paper. The author also thankful to the referees for providing valuable times for the article.

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Corrigendum

Certain typographical mistakes are communicated in the paper of Sengupta [4]:

(i) On page 385: The equation $y_i(x) = \dots$ in the third line will be numbered (15).

(ii) On page 396: In the last line:

Choose $\phi^*(x, \lambda_n) = b_{2n}U_1(x, \lambda_n) - b_{1n}(x, \lambda_n)\overline{U_3}(x, \lambda_n)$ an eigenvector.