CONTINUOUS FOURTH DERIVATIVE METHOD
FOR THIRD ORDER BOUNDARY VALUE PROBLEMS

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Abstract: A fourth derivative method (FDM) with continuous coefficients is derived and used to obtain main and additional methods which are used to solve third order boundary value problems (TOBVPs) such as the Blasius, the Falkner-Skan, and the sandwich beam problems which frequently occur in engineering. The convergence analysis of the method is discussed. Numerical experiments are performed to show speed and accuracy advantages.

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1. Introduction

In the last three decades, third order boundary value problems (BVP’s) have gained a lot of attention in the literature. These kind of BVP’s have applications in the field of engineering and science, especially in control theory, and biological sciences. For instance, draining and coating flow problems, see Tuck and Schwartz [21]; laminar boundary layer and sandwich beam problems, see
In the past decades, tremendous attention has been focused on developing methods for the solution of \( y''' = f(x, y, y', y'') \) subject to boundary conditions (see Awoyemi [2], Jator [6], [7]). Many methods like standard finite difference (SFDs), splines (Khan and Aziz [12]), non-polynomial splines (Siraj-ul-Islam and Tirmizi [19]), high order difference methods (Salama and Mansour [17]) are used in solving third order BVP’s. Most of these methods are solved by first reducing a higher order ordinary differential equation (ODE) to an equivalent system of first-order ODE’s which takes a lot of human effort and computer time as discussed in Awoyemi [2].

In this paper, we consider the general third-order BVP

\[
Dy \equiv y''' = f(x, y, y', y'')
\]  

subject to any of the following boundary conditions:

\[
\begin{align*}
y(a) &= y_0, \quad y'(a) = \delta_0, \quad y(b) = y_N \\
y(a) &= y_0, \quad y'(a) = \delta_0, \quad y(b) = y_M \\
y(a) &= y_0, \quad y'(b) = y_M, \quad y(c) = \gamma_1, \text{ where } c = \left(\frac{a+b}{2}\right)
\end{align*}
\]

where \( a, b, c, \delta_0, y_0, y_N, y_M \) are constants, \( f \) is a continuous function and satisfies a Lipschitz condition as given in [4]. Keller [10] has given the theorem and the proof of the general conditions which ensure that the solution to (1) will exist and be unique.

Using multistep collocation a continuous FDM is derived, see Lie and Norsett [13], Atkinson [1], and Onumanyi et al [15]. The continuous representation generates the basic FDM and \( k-1 \) additional methods which are combined and used to simultaneously produce approximations \( y_j \), for \( j = 1, \ldots, N-1 \) to the solution of (1) at points \( x_j \) for \( j = 1, \ldots, N-1 \) on \( \pi_N : a = x_0 < x_1 < \ldots < x_N = b \)

The basic and auxiliary methods are obtained from the same continuous scheme and are of the same order, hence, possible errors which are due to auxiliary methods of lower order are avoided as the integration proceeds on the entire interval.

The paper is organized as follows. In section two, we derive an approximation \( U(x) \) for \( y(x) \) which is used to obtain the main and additional FDMs. The convergence analysis of the method is also given in Section 2. Section 3 is devoted to the computational aspects and an algorithm equipped with an automatic error estimate based on the double mesh principle. Numerical examples are given in Section 4 to show speed and accuracy advantages. Finally, the conclusion of the paper is discussed in Section 5.
2. Construction of the CFDMs with FDMs as by Products

In this section, we construct the CFDMs with FDMs as by products. Thus, on the interval \([x_n, x_{n+k}]\), we assume that the exact solution \(y(x)\) and its first and second derivatives are locally represented by the continuous methods

\[
U_k(x) = \sum_{j=0}^{k-1} \alpha_j(x)y_{n+j} + h^3 \sum_{j=0}^{k} \beta_j(x)f_{n+j} + h^4 \sum_{j=0}^{k} \gamma_j(x)g_{n+j} \quad (2)
\]

\[
\begin{align*}
U_k'(x) &= \frac{d}{dx}(U(x)) \\
U_k''(x) &= \frac{d^2}{dx^2}(U(x))
\end{align*}
\quad (3)
\]

where \(\alpha_j(x), \beta_j(x), \gamma_j(x)\) are continuous coefficients, and \(m > 0\) is an integer.

We assume that \(y_{n+j} = U_k(x_n + jh)\) is the numerical approximation to the analytical solution \(y(x_{n+j})\), \(y_{n+j} = U_k(x_n + jh)\) is an approximation to \(y'(x_{n+j})\), \(y''_{n+j} = U_k''(x_n + jh)\) is an approximation to \(y''(x_{n+j})\), \(f_{n+j} = U_k'''(x_n + jh)\), and \(g_{n+j} = U_k''''(x_n + jh)\). We also note that \(f_{n+j} = f(x_{n+j}, y_{n+j}, y'_{n+j}, y''_{n+j}), g_{n+j} = \frac{df(x,y(x),y'(x),y''(x))}{dx}x_{n+j}, y_{n+j}, y'_{n+j}, y''_{n+j}, j = 0, \ldots, 3\).

The continuous method (2) and its first and second derivatives (3) are piece-wise continuous on \([a, b]\) and defined for all \(x \epsilon [a, b]\). That is, \(U_k(x), U'_k(x), U''_k(x)\) are defined such that \(U_k(x) = y(x) + O(h^{11}), U'_k(x) = \frac{d}{dx}(y(x) + O(h^{11})), U''_k(x) = \frac{d^2}{dx^2}(y(x) + O(h^{11})), x \epsilon (x_n, x_{n+3})\). The polynomials \(\{U_0(x), U_3(x), \ldots, U_{N-3}(x)\}, \{U_0(x), U_3(x), \ldots, U_{N-3}(x)\}, \{U_0''(x), U_3''(x), \ldots, U_{N-3}''(x)\}, \{U_0''(x), U_3''(x), \ldots, U_{N-3}''(x)\}\), then define piece-wise polynomials \(U(x), U'(x), \) and \(U''(x)\) which are also continuous on \([a, b]\).

Hence, (2) and (3) have the ability to provide a continuous solution on \([a, b]\) with a uniform accuracy comparable to that obtained at the grid points (see Onumanyi et al. [15]) and can also be used to produce additional discrete methods (see Onumanyi et al. [14]). In what follows, we state the theorem that facilitates the construction of the CFDMs (2) and (3).

**Theorem 2.1.** Let the following conditions be satisfied

\[
U_k(x_{n+j}) = y_{n+j}, j = 0, 1, 2 \quad (4)
\]

\[
U_k'''(x_{n+j}) = f_{n+j}, U_k''''(x_{n+j}) = g_{n+j}, j = 0, \ldots, 3, \quad (5)
\]

then, the continuous representations (2) and (3) are equivalent to the following:

\[
U_k(x) = V^T \left(W^{-1}\right)^T P(x) \quad (6)
\]
\[
\begin{align*}
U'_k(x) &= \frac{d}{dx}(V^T (W^{-1})^T P(x)) \\
U''_k(x) &= \frac{d^2}{dx^2}(V^T (W^{-1})^T P(x))
\end{align*}
\]  

(7)

where \(W\) is given as

\[
W = \begin{pmatrix}
P_0(x_n) & \cdots & P_{10}(x_n) \\
P_0(x_{n+1}) & \cdots & P_{10}(x_{n+1}) \\
P_0(x_{n+2}) & \cdots & P_{10}(x_{n+2}) \\
P'_0(x_n) & \cdots & P'_{10}(x_n) \\
P'_0(x_{n+1}) & \cdots & P'_{10}(x_{n+1}) \\
P'_0(x_{n+2}) & \cdots & P'_{10}(x_{n+2}) \\
P''_0(x_n) & \cdots & P''_{10}(x_n) \\
P''_0(x_{n+1}) & \cdots & P''_{10}(x_{n+1}) \\
P''_0(x_{n+2}) & \cdots & P''_{10}(x_{n+2}) \\
P'''_0(x_n) & \cdots & P'''_{10}(x_n)
\end{pmatrix},
\]

\[V = (y_n, y_{n+1}, y_{n+2}, f_n, f_{n+1}, f_{n+2}, f_{n+3}, g_n, g_{n+1}, g_{n+2}, g_{n+3})^T,\]

\[P(x) = (P_0(x), P_1(x), \ldots, P_{10}(x))^T.\]

We note that \(T\) denotes the transpose and \(P_j(x) = x^j, j = 0, \ldots, 10\) are basis functions.

\textbf{Proof.} The proof is the same as in Jator et al. [9].

\[\square\]

\textbf{Remark 2.2.} We emphasize that the continuous methods (2) and (3) which are equivalent to the forms (6) and (7) are used to produce the main and additional methods which are combined and simultaneously applied to provide all approximations on the entire interval for boundary value problems of the form \(y''' = f(x, y, y', y'').\)

\textbf{Remark 2.3.} The continuous methods (2) and (3) are obtained by solving a system of 11 equations resulting from conditions (4) and (5) given in theorem 2.1.

\textbf{Coefficients of the CFDMs (2).} In order to simplify the coefficients of (2), we introduce the scale variable \(t = \frac{x - x_{n+2}}{h}\) to express the coefficients of the CFDM as follows:

\[\alpha_0(t) = \frac{1}{2}(t + t^2) \quad ; \quad \alpha_1(t) = (-2t - t^2) \quad ; \quad \alpha_2(t) = \frac{1}{2}(2 + 3t + t^2)\]
\beta_0(t) = \frac{1}{544320} (6950t + 7891t^2 + 2100t^5 + 462t^6 - 1200t^7 - 330t^8 + 250t^9 + 77t^{10})
\beta_1(t) = \frac{1}{20160} (4268t + 4886t^2 + 672t^5 - 168t^6 - 240t^7 + 15t^8 + 40t^9 + 7t^{10})
\beta_2(t) = \frac{1}{20160} (2094t + 4759t^2 + 3360t^3 - 924t^5 - 42t^6 + 240t^7 + 30t^8 - 30t^9 - 7t^{10})
\beta_3(t) = \frac{1}{544320} (2716t + 3854t^2 + 4704t^5 + 5208t^6 + 1200t^7 - 885t^8 - 520t^9 - 7t^{10})
\gamma_0(t) = \frac{1}{181440} (430t + 503t^2 + 168t^5 + 42t^6 - 96t^7 - 30t^8 + 20t^9 + 7t^{10})
\gamma_1(t) = + \frac{1}{20160} (228t + 434t^2 + 336t^5 - 168t^7 - 15t^8 + 30t^9 + 7t^{10})
\gamma_2(t) = \frac{1}{20160} (-610t - 1009t^2 + 840t^4 + 336t^5 - 294t^6 - 192t^7 - 30t^8 + 40t^9 + 7t^{10})
\gamma_3(t) = + \frac{1}{181440} (-212t - 298t^2 - 336t^5 - 336t^6 - 24t^7 + 105t^8 + 50t^9 + 7t^{10})

\textbf{FDMs.} The following main methods are obtained by evaluating (2) and (3) at } x = x_{n+3}.

\begin{align*}
y_{n+3} - 3y_{n+2} + 3y_{n+1} - y_n &= \frac{h^3}{168} (5f_n + 79f_{n+1} + 79f_{n+2} + 5f_{n+3}) \\
&\quad + \frac{h^4}{5040} (29g_n + 213g_{n+1} - 213g_{n+2} - 29g_{n+3}), \quad (8)
\end{align*}

\begin{align*}
h^2y'_{n+3} - \frac{3}{2}y_n + 4y_{n+1} - \frac{5}{2}y_{n+2} &= \frac{h^3}{90720} (4664f_n + 68679f_{n+1} + 82800f_{n+2} + 10177f_{n+3}) \\
&\quad + \frac{h^4}{30240} (311g_n + 2730g_{n+1} - 39g_{n+2} - 552g_{n+3}), \quad (9)
\end{align*}

\begin{align*}
h^2y''_{n+3} - y_n + 2y_{n+1} - y_{n+2} &= \frac{h^3}{272160} (13846f_n + 167967f_{n+1} + 255258f_{n+2} \\
&\quad + 107249f_{n+3}) + \frac{h^4}{90720} (992g_n + 11169g_{n+1} + 13896g_{n+2} - 4147g_{n+3}), \quad (10)
\end{align*}

The following additional methods are obtained by evaluating (3) at } x =
\[ -2y_1 = -h y_0 - \frac{3}{2} y_0 - \frac{1}{2} y_2 + \frac{h^3}{30240} (6127f_0 + 18810f_1 + 4689f_2 + 614f_3) \\
+ \frac{h^4}{30240} (291g_0 - 1878g_1 - 813g_2 - 50g_3) \\
- y_2 = -h^2 y_0 + y_0 - 2y_1 + \frac{h^3}{272160} (-99149f_0 - 127278f_1 - 39987f_2 - 5746f_3) \\
+ \frac{h^4}{18144} (-725g_0 + 3546g_1 + 1467g_2 + 94g_3) \\
h y_1' = -\frac{1}{2} y_0 + \frac{1}{2} y_2 + \frac{h^3}{272160} (-2692f_0 - 35721f_1 - 6372f_2 - 575f_3) \\
+ \frac{h^4}{90720} (-155g_0 + 828g_1 + 891g_2 + 46g_3) \\
h y_2' = \frac{1}{2} y_0 - 2y_1 + \frac{3}{2} y_2 + \frac{h^3}{272160} (3475f_0 + 57618f_1 + 28269f_2 + 1358f_3) \\
+ \frac{h^4}{90720} (215g_0 + 1026g_1 - 2745g_2 - 106g_3) \\
h^2 y_1'' = y_0 - 2y_1 + y_2 + \frac{h^3}{272160} (4246f_0 - 513f_1 - 3942f_2 + 209f_3) \\
+ \frac{h^4}{90720} (224g_0 - 5247g_1 + 72g_2 - 19g_3) \\
h^2 y_2'' = y_0 - 2y_1 + y_2 + \frac{h^3}{272160} (7891f_0 + 131922f_1 + 128493f_2 + 3854f_3) \\
+ \frac{h^4}{90720} (503g_0 + 3906g_1 - 9081g_2 - 298g_3) \\
\] 

(11)

3. Convergence Analysis

**Local truncation error and order.** The local truncation errors associated with the discrete fourth derivative method (FDM) main methods (8)-(10) are given by

\[
\tau_1 = \frac{47}{84006000} h^{11} y^{(11)}(x_i + \theta_i) + O(h^{12}), \quad h \tau_1' = \frac{6641}{558835200} h^{11} y^{(11)}(x_i + \theta_i) + O(h^{12}),
\]

\[
h^2 \tau_1'' = \frac{29}{16934400} h^{11} y^{(11)}(x_i + \theta_i) + O(h^{12}), \quad i = 3, \ldots, N, \quad |\theta_i| \leq 1.
\]

The local truncation errors associated with the first, second and third additional methods in (11) are given by

\[
\tau_1 = \frac{71}{19958400} h^{11} y^{(11)}(\xi) + O(h^{12}), \quad h \tau_1 = \frac{59}{62092800} h^{11} y^{(11)}(\xi) + O(h^{12}), \quad h^2 \tau_1'' = \frac{17}{9313920} h^{11} y^{(11)}(\xi) + O(h^{12}), \quad x_0 \leq \xi \leq x_1.
\]

Similarly, the local truncation errors associated with the fourth, fifth and sixth additional methods in (11) are given by

\[
\tau_2 = \frac{1}{86400} h^{11} y^{(11)}(\xi) + O(h^{12}), \quad h \tau_2 = \frac{19}{25401600} h^{11} y^{(11)}(\xi) + O(h^{12}), \quad h^2 \tau_2'' = \frac{61}{12700800} h^{11} y^{(11)}(\xi) + O(h^{12}), \quad x_1 \leq \xi \leq x_2.
\]

**Convergence.** In order to show that the FDM converges, the methods (8), (9), (10), and (11) can be compactly written in matrix form by introducing the
following notations. Let $A$ be a $3N \times 3N$ matrix defined by

$$A = \begin{bmatrix} A_{11} & A_{12} & A_{13} \\ A_{21} & A_{22} & A_{23} \\ A_{31} & A_{32} & A_{33} \end{bmatrix},$$

where the elements of $A$ are $N \times N$ matrices given as

$$A_{11} = \begin{bmatrix} -2 & \frac{1}{2} & 0 & 0 & 0 & 0 & \cdots & 0 & 0 \\ 2 & -1 & 0 & 0 & 0 & 0 & \cdots & 0 & 0 \\ 3 & -3 & 1 & 0 & 0 & 0 & \cdots & 0 & 0 \\ -1 & 3 & -3 & 1 & 0 & 0 & \cdots & 0 & 0 \\ \vdots & \ddots & \vdots & \vdots & \vdots & \vdots & \cdots & \vdots & \vdots \\ 0 & 0 & 0 & \cdots & -1 & 3 & -3 & 1 & 0 \\ 0 & 0 & 0 & \cdots & 0 & -1 & 3 & -3 & 1 \end{bmatrix},$$

$$A_{31} = \begin{bmatrix} 2 & -1 & 0 & 0 & 0 & 0 & \cdots & 0 & 0 \\ 2 & -1 & 0 & 0 & 0 & 0 & \cdots & 0 & 0 \\ 2 & -1 & 0 & 0 & 0 & 0 & \cdots & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 & 0 & \cdots & 0 & 0 \\ \vdots & \ddots & \vdots & \vdots & \vdots & \vdots & \cdots & \vdots & \vdots \\ 0 & 0 & \cdots & -1 & 2 & -1 & 0 & 0 \\ 0 & 0 & \cdots & 0 & -1 & 2 & -1 & 0 \end{bmatrix},$$

$$A_{21} = \begin{bmatrix} 0 & -1/2 & 0 & 0 & 0 & 0 & \cdots & 0 & 0 \\ 2 & -3/2 & 0 & 0 & 0 & 0 & \cdots & 0 & 0 \\ 4 & -5/2 & 0 & 0 & 0 & 0 & \cdots & 0 & 0 \\ -3/2 & 4 & -5/2 & 0 & 0 & 0 & \cdots & 0 & 0 \\ \vdots & \ddots & \vdots & \vdots & \vdots & \vdots & \cdots & \vdots & \vdots \\ 0 & 0 & \cdots & -3/2 & 4 & -5/2 & 0 \\ 0 & 0 & \cdots & 0 & -3/2 & 4 & -5/2 \end{bmatrix},$$

$$A_{12} = A_{13} = A_{32} = \begin{bmatrix} 0 & 0 & 0 & \cdots & 0 \\ 0 & 0 & 0 & \cdots & 0 \\ 0 & 0 & 0 & \cdots & 0 \\ \vdots & \ddots & \vdots & \cdots & \vdots \\ 0 & 0 & 0 & \cdots & 0 \end{bmatrix},$$
Similarly, let $B$ be a $3N \times 3N$ matrix defined by

$$B = \begin{bmatrix} B_{11} & B_{12} & B_{13} \\ B_{21} & B_{22} & B_{23} \\ B_{31} & B_{32} & B_{33} \end{bmatrix},$$

where the elements of $B$ are $N \times N$ matrices given as

$$B_{11} = h^3 \begin{bmatrix} 18810 & 4689 & 614 \\ -127278 & 90720 & -5746 \\ 168 & 168 & 168 \end{bmatrix},$$

$$B_{12} = h^3 \begin{bmatrix} -1878 & -813 & -50 \\ 30240 & 30240 & 30240 \\ 5040 & 5040 & 5040 \end{bmatrix},$$

$$B_{21} = h^3 \begin{bmatrix} -35721 & -6372 & -575 \\ 272160 & 272160 & 272160 \\ 90720 & 90720 & 90720 \end{bmatrix},$$

$$A_{22} = A_{33} = \begin{bmatrix} 1 & 0 & \ldots & 0 \\ 0 & 1 & 0 & \ldots & 0 \\ 0 & 0 & 1 & \ldots & 0 \\ \vdots & \ddots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \ldots & 1 \end{bmatrix}.$$
We further define the following vectors:

\[ Y = (y(x_1), \ldots, y(x_N), h y'(x_1), \ldots, h y'(x_N), h^2 y''(x_1), \ldots, h^2 y''(x_N))^T, \]

\[ \bar{Y} = (y_1, \ldots, y_N, h y_1', \ldots, h y_N', h^2 y_1'', \ldots, h^2 y_N'')^T, \]

\[ F = (f_1, \ldots, f_N, h g_1, \ldots, h g_N, h^2 z_1, \ldots, h^2 z_N)^T \]
(the variables $z_i$ are introduced with coefficients zeros to augment the matrix $B$),

$$C = \left( hy_0 - \frac{3}{2} y_0 - \frac{6127}{90720} h^3 f_0 - \frac{291}{30240} h^4 g_0, h^2 y_0 - y_0 + \frac{99149}{272160} h^3 f_0 \right)$$

$$+ \left( \frac{725}{18144} h^4 g_0, -y_0 + \frac{5}{168} h^3 f_0 - \frac{29}{5040} h^4 g_0, 0, \ldots, 0, \frac{1}{2} y_0 + \frac{2692}{272160} h^3 f_0 \right)$$

$$+ \left( \frac{155}{90720} h^4 g_0, -\frac{1}{2} y_0 - \frac{3475}{272160} h^3 f_0 - \frac{215}{90720} h^4 g_0, -\frac{3}{2} y_0 - \frac{4644}{90720} h^3 f_0 \right)$$

$$- \left( \frac{311}{30240} h^4 g_0, 0, \ldots, 0, -y_0 - \frac{4246}{272160} h^3 f_0 - \frac{224}{90720} h^4 g_0, -y_0 - \frac{7891}{272160} h^3 f_0 \right)$$

$$- \left( \frac{503}{90720} h^4 g_0, 0, \ldots, 0, \frac{13846}{90720} h^3 f_0 - \frac{922}{30240} h^4 g_0 \right)$$

$$L(h) = (\tau_1, \ldots, \tau_N, h\tau_1, \ldots, h\tau_N, h^2\tau_1, \ldots, h^2\tau_N)^T,$$

where $L(h)$ is the local truncation error.

$$E = \bar{Y} - Y = (e_1, \ldots, e_N, he_1, \ldots, he_N, h^2e_1, \ldots, h^2e_N)^T.$$

**Theorem 3.1.** Let $\bar{Y}$, $Y$, and $E$ be defined as above. Let $\bar{Y}$ be an approximation of the solution vector $Y$ for the system formed by combining the methods (8), (9), (10), and (11). If $e_i = |y(x_i) - y_i|$, $he_i = |hy(x_i) - hy_i|$, and $h^2e_i = |h^2y''(x_i) - h^2y''_i|$ where the exact solution $y(x)$ is several times differentiable on $[a, b]$ and if $\|E\| = \|Y - \bar{Y}\|$, then, the FDM is an eighth-order convergent method. That is $\|E\| = O(h^8)$.

**Proof.** The exact form of the system is given by (12)

$$AY - BF(Y) + C + L(h) = 0, \quad (12)$$

and the approximate form of the system is given by (13)

$$A\bar{Y} - BF(\bar{Y}) + C = 0, \quad (13)$$

where $\bar{Y}$ is the approximation of the solution vector $Y$.

Subtracting (12) from (13) we obtain

$$AE - BF(\bar{Y}) + BF(Y) = L(h), \quad (14)$$

Using the mean-value theorem, we write (14) as

$$(A - BJ)E = L(h),$$
where the Jacobian matrix $J$ and its entries $J_{11}, J_{12}, J_{13}, J_{21}, J_{22}, J_{23}, J_{31}, J_{32},$ and $J_{33}$ are defined as follows:

$$J = \begin{bmatrix} J_{11} & J_{12} & J_{13} \\ J_{21} & J_{22} & J_{23} \\ J_{31} & J_{32} & J_{33} \end{bmatrix}, \quad J_{11} = \begin{bmatrix} \frac{\partial f_1}{\partial y_1} & \cdots & \frac{\partial f_1}{\partial y_N} \\ \frac{\partial f_2}{\partial y_1} & \cdots & \frac{\partial f_2}{\partial y_N} \\ \cdots & \cdots & \cdots \end{bmatrix},$$

$$J_{12} = \begin{bmatrix} \frac{\partial f_1}{\partial y_1} & \cdots & \frac{\partial f_1}{\partial y_N} \\ \frac{\partial f_2}{\partial y_1} & \cdots & \frac{\partial f_2}{\partial y_N} \\ \cdots & \cdots & \cdots \end{bmatrix}, \quad J_{13} = \begin{bmatrix} \frac{\partial f_1}{\partial y_1} & \cdots & \frac{\partial f_1}{\partial y_N} \\ \frac{\partial f_2}{\partial y_1} & \cdots & \frac{\partial f_2}{\partial y_N} \\ \cdots & \cdots & \cdots \end{bmatrix},$$

$$J_{21} = h \begin{bmatrix} \frac{\partial g_1}{\partial y_1} & \cdots & \frac{\partial g_1}{\partial y_N} \\ \frac{\partial g_2}{\partial y_1} & \cdots & \frac{\partial g_2}{\partial y_N} \\ \cdots & \cdots & \cdots \end{bmatrix}, \quad J_{22} = h \begin{bmatrix} \frac{\partial g_1}{\partial y_1} & \cdots & \frac{\partial g_1}{\partial y_N} \\ \frac{\partial g_2}{\partial y_1} & \cdots & \frac{\partial g_2}{\partial y_N} \\ \cdots & \cdots & \cdots \end{bmatrix},$$

$$J_{23} = h \begin{bmatrix} \frac{\partial g_1}{\partial y_1} & \cdots & \frac{\partial g_1}{\partial y_N} \\ \frac{\partial g_2}{\partial y_1} & \cdots & \frac{\partial g_2}{\partial y_N} \\ \cdots & \cdots & \cdots \end{bmatrix}, \quad J_{31} = h^2 \begin{bmatrix} \frac{\partial z_1}{\partial y_1} & \cdots & \frac{\partial z_1}{\partial y_N} \\ \frac{\partial z_2}{\partial y_1} & \cdots & \frac{\partial z_2}{\partial y_N} \\ \cdots & \cdots & \cdots \end{bmatrix},$$

$$J_{32} = h^2 \begin{bmatrix} \frac{\partial z_1}{\partial y_1} & \cdots & \frac{\partial z_1}{\partial y_N} \\ \frac{\partial z_2}{\partial y_1} & \cdots & \frac{\partial z_2}{\partial y_N} \\ \cdots & \cdots & \cdots \end{bmatrix}, \quad J_{33} = h^2 \begin{bmatrix} \frac{\partial z_1}{\partial y_1} & \cdots & \frac{\partial z_1}{\partial y_N} \\ \frac{\partial z_2}{\partial y_1} & \cdots & \frac{\partial z_2}{\partial y_N} \\ \cdots & \cdots & \cdots \end{bmatrix}. $$

Let $M = -BJ$ be a matrix of dimension $3N$. We have

$$(A + M)E = L(h), \quad (15)$$

and for sufficiently small $h$, $A + M$ is a monotone matrix and thus invertible (see Jain and Aziz [5] and Jator and Li [8]). Therefore,

$$(A + M)^{-1} = D = (d_{i,j}) \geq 0, \quad \text{and} \quad \sum_{j=1}^{3N} d_{i,j} = O(h^{-3}). \quad (16)$$

If $\|E\|$ is the norm of maximum global error and from (15), $E = (A + M)^{-1}L(h)$, using (16) and the truncation error vector $L(h)$, it follows that

$$\|E\| = O(h^{8}).$$

Therefore, the TDM is an eighth-order convergent method. □
4. Numerical Examples

In this section, we test three linear and two non-linear numerical examples to illustrate the accuracy of the method. The global maximum absolute error is computed as $E_{\text{MAX}} = \text{Max} |y(x_i) - y_i|$, $i = 1, \ldots, N$, where $y(x_i)$ is the exact solution computed at the grid point and $y_i$ is an approximation to the exact solution using the FDM. We note that the method requires only two function evaluations per step. All computations were carried out using our written code in Mathematica 8.0.

**Example 1.** Consider the linear third order BVP on $0 \leq x \leq 1$.

$$y''' - xy = (x^3 - 2x^2 - 5x - 3)e^x,$$

$y(0) = y(1) = 1$, and $y'(0) = 1$

The exact solution of the system is given by

$$y(x) = x(1 - x)e^x$$

The FDM was tested on example 1 and results were compared with Brugnano and Trigiante (BT) (order $p = 8$). As expected, the FDM performed better than BT.

<table>
<thead>
<tr>
<th>$N$</th>
<th>FDM $E_{\text{MAX}}$</th>
<th>BT [3] $E_{\text{MAX}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>7</td>
<td>$4.12 \times 10^{-12}$</td>
<td>$2.89 \times 10^{-8}$</td>
</tr>
<tr>
<td>14</td>
<td>$1.56 \times 10^{-14}$</td>
<td>$6.93 \times 10^{-11}$</td>
</tr>
<tr>
<td>28</td>
<td>$6.08 \times 10^{-17}$</td>
<td>$1.28 \times 10^{-13}$</td>
</tr>
<tr>
<td>56</td>
<td>$2.37 \times 10^{-19}$</td>
<td>$2.69 \times 10^{-15}$</td>
</tr>
<tr>
<td>112</td>
<td>$9.27 \times 10^{-22}$</td>
<td>$1.448 \times 10^{-15}$</td>
</tr>
</tbody>
</table>

Table 1: Results for Example 1

**Example 2.** Consider the sandwich beam problem (linear third order) BVP on $0 \leq x \leq 1$ found by Krajcinovic [11].

$$y''' - l^2 y' + a = 0, \quad y'(0) = y'(1) = y\left(\frac{1}{2}\right) = 0$$

The exact solution of the system is given by
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\[ y(x) = \left(\frac{a}{l^3}\right) \left[ \left( \sinh \left( \frac{l}{2} \right) - \sinh(lx) \right) + l \left( x - \frac{1}{2} \right) + \tanh \left( \frac{l}{2} \right) \left( \cosh(lx) - \cosh \left( \frac{l}{2} \right) \right) \right] \]

We tested the method for \( a=1, l=5, \) and \( 10, \) and the results are compared with Brugnano and Trigiante (BT)(order \( p = 8 \)). It is obvious from the numerical results in Table 2 that our method performed excellently when compared with Brugnano and Trigiante method.

<table>
<thead>
<tr>
<th>( l = 5 )</th>
<th>FDM</th>
<th>BT</th>
<th>( l = 10 )</th>
<th>FDM</th>
<th>BT</th>
</tr>
</thead>
<tbody>
<tr>
<td>( N )</td>
<td>( 14 )</td>
<td>( 28 )</td>
<td>( 56 )</td>
<td>( 112 )</td>
<td>( 14 )</td>
</tr>
<tr>
<td>( 5.78 \times 10^{-12} )</td>
<td>( 2.16 \times 10^{-14} )</td>
<td>( 4.94 \times 10^{-16} )</td>
<td>( 1.07 \times 10^{-16} )</td>
<td>( 1.47 \times 10^{-9} )</td>
<td>( 6.70 \times 10^{-12} )</td>
</tr>
</tbody>
</table>

Table 2: Results for Example 2

**Example 3.** Consider the non-linear third order BVP on \( 0 \leq x \leq 1 \).

\[ y''' = -2e^{-3y} + 4(1 + x)^{-3}, \quad y(0) = 0, \quad y'(0) = 1, \quad y(1) = \ln 2 \]

The exact solution of the system is given by

\[ y(x) = \ln(1 + x) \]

Clearly, our results are better than those of Brugnano and Trigiante [3]. Details of the numerical results are given in Table 3.

<table>
<thead>
<tr>
<th>( N )</th>
<th>FDM</th>
<th>BT</th>
</tr>
</thead>
<tbody>
<tr>
<td>( 7 )</td>
<td>( 5.24 \times 10^{-9} )</td>
<td>( 4.13 \times 10^{-7} )</td>
</tr>
<tr>
<td>( 14 )</td>
<td>( 2.39 \times 10^{-11} )</td>
<td>( 5.53 \times 10^{-9} )</td>
</tr>
<tr>
<td>( 28 )</td>
<td>( 9.50 \times 10^{-14} )</td>
<td>( 4.63 \times 10^{-11} )</td>
</tr>
<tr>
<td>( 56 )</td>
<td>( 3.62 \times 10^{-16} )</td>
<td>( 2.16 \times 10^{-13} )</td>
</tr>
<tr>
<td>( 112 )</td>
<td>( 2.27 \times 10^{-17} )</td>
<td>( 8.83 \times 10^{-16} )</td>
</tr>
</tbody>
</table>

Table 3: Results for Example 3

**Example 4.** Consider the boundary-layer problem (nonlinear third order) taken from Salama [16].
We will consider the Falkner-Skan Equation ($\alpha = 1$) and the Blasius equations ($\alpha = \frac{1}{2}$ and $\beta = 0$) that are special cases of the above boundary layer problem. The results to the above solution are presented differently as there are no theoretical solutions. We looked at different values of $\beta$ (positive, zero and negative). The numerical results are shown in Table 4 for the Falkner-Skan Equation and in Table 5 for the Blasius Equation. It is clear from the table that our results are more efficient. The Number of steps needed in our method was only 10 to get the required values at the truncated boundary, where as in Brugnano and Trigiante [3], the number of steps needed was 21, and in Salama [16] ($\varepsilon = 10^{-2}$), the number of steps needed was 200.

<table>
<thead>
<tr>
<th>$\eta$</th>
<th>$\beta$</th>
<th>$y''(0)$</th>
<th>$y'(0)$</th>
<th>$y''(0)$</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.93</td>
<td>40</td>
<td>7.314787</td>
<td>7.3149</td>
<td>7.314787</td>
</tr>
<tr>
<td>0.95</td>
<td>30</td>
<td>6.338219</td>
<td>6.33826</td>
<td>6.338219</td>
</tr>
<tr>
<td>1.13</td>
<td>20</td>
<td>5.180731</td>
<td>5.18076</td>
<td>5.180731</td>
</tr>
<tr>
<td>1.49</td>
<td>10</td>
<td>3.675257</td>
<td>3.67527</td>
<td>3.675257</td>
</tr>
<tr>
<td>2.57</td>
<td>2</td>
<td>1.687317</td>
<td>1.68732</td>
<td>1.687317</td>
</tr>
<tr>
<td>2.88</td>
<td>1</td>
<td>1.232951</td>
<td>1.23295</td>
<td>1.2329519</td>
</tr>
<tr>
<td>3.29</td>
<td>0.5</td>
<td>0.928234</td>
<td>0.928234</td>
<td>0.928234</td>
</tr>
<tr>
<td>4.01</td>
<td>0</td>
<td>0.47110</td>
<td>0.471107</td>
<td>0.47110</td>
</tr>
<tr>
<td>4.27</td>
<td>-0.1</td>
<td>0.321838</td>
<td>0.321838</td>
<td>0.321838</td>
</tr>
<tr>
<td>4.49</td>
<td>-0.15</td>
<td>0.220244</td>
<td>0.220245</td>
<td>0.220244</td>
</tr>
<tr>
<td>4.71</td>
<td>-0.18</td>
<td>0.134948</td>
<td>0.134941</td>
<td>0.134948</td>
</tr>
<tr>
<td>5.00</td>
<td>-0.1988</td>
<td>0.039868</td>
<td>0.039859</td>
<td>0.039868</td>
</tr>
</tbody>
</table>

Table 4: Results for Falkner-Skan Equation for different values of $\beta$

The results of the Blasius Equation are given below:

5. Conclusions

We have derived a CFDM from which discrete FDMs are obtained and applied to solve $y''' = f(x, y, y', y'')$ subject to Dirichlet and Neumann boundary conditions without first adapting the ODE to an equivalent first order system. Numerical experiments are performed that show efficiency and accuracy advan-
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Table 5: Results for different values of the truncated boundary $\eta_\infty$ for Blasius Equation

tages of the method over existing ones in the literature. Details of the numerical results are displayed in Tables 1, 2, 3, 4, 5, 6.

References


