

**A NOTE ON THE HYPERCOMPLEX RIEMANN-CAUCHY  
LIKE RELATIONS FOR QUATERNIONS  
AND LAPLACE EQUATIONS**

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**Abstract:** In this Note it is worked out a new set of Laplace-Like equations for quaternions through Riemann-Cauchy hypercomplex relations obtained earlier [1]. As in the theory of functions of a complex variable, it is expected that this new set of Laplace-Like equations might be applied to a large number of Physical problems, providing new insights in the Classical Fields Theory.

**AMS Subject Classification:** 30G99, 30E99

**Key Words:** quaternions, Laplace's equations, Riemann-Cauchy relations

**1. Cauchy-Riemann Equations (Functions of One Complex Variable)**

In order to fix ideas will be considered theorem that relates the partial derivatives for the case of a function  $f(z)$  of a complex variable  $f(z) = u(x, y) + iv(x, y)$  [2], which here will be called Riemann-Cauchy conditions. These relations say that the first order partial derivatives of functions  $u(x, y)$  and  $v(x, y)$  satisfy relations according to the following theorem:

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Received: May 20, 2013

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**Theorem 1.** *Is  $f(z) = u(x, y) + iv(x, y)$  a function defined and continues in a neighborhood of the point  $z = x + yi$  and differentiable at  $z$ . Then the partial derivatives of the first order of  $u(x, y)$  and  $v(x, y)$  exist and satisfy the relations:*

$$\frac{\partial u(x, y)}{\partial x} = \frac{\partial v(x, y)}{\partial y}, \quad (1)$$

$$\frac{\partial u(x, y)}{\partial y} = -\frac{\partial v(x, y)}{\partial x}. \quad (2)$$

Thus, if  $f(z)$  is analytic in a domain  $\Gamma$ , its partial derivatives exist and satisfy the set of relations (1) and (2) over all point in  $\Gamma$ . Moreover, with the above functions class  $C^2$  using Schwartz's Theorem for partial derivatives immediately follows the following equations:

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0, \quad (3)$$

$$\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} = 0. \quad (4)$$

the above equations are called Laplace's equations.

## 2. Cauchy-Riemann Conditions for Quaternionic Functions

Now we may consider a set of conditions presented as the Riemann-Cauchy like relations for quaternionic functions. It follows the Theorem [1]:

**Theorem 2.** *For any pair points  $a$  and  $b$  and any path joining them simply connect subdomain of the four-dimensional space, the integral  $\int_a^b f dq$  is independent from the given path if and only if there is a function  $F = F_1 + F_2i + F_3j + F_4k$  such that  $\int_a^b f dq = F(a) - F(b)$ , and satisfying the following relations:*

$$\frac{\partial F_1}{\partial x_1} = \frac{\partial F_2}{\partial x_2} = \frac{\partial F_3}{\partial x_3} = \frac{\partial F_4}{\partial x_4}, \quad (5)$$

$$\frac{\partial F_2}{\partial x_1} = -\frac{\partial F_1}{\partial x_2} = -\frac{\partial F_3}{\partial x_4} = \frac{\partial F_4}{\partial x_3}, \quad (6)$$

$$\frac{\partial F_3}{\partial x_1} = -\frac{\partial F_1}{\partial x_3} = -\frac{\partial F_2}{\partial x_4} = \frac{\partial F_4}{\partial x_2}, \quad (7)$$

$$\frac{\partial F_4}{\partial x_1} = \frac{\partial F_1}{\partial x_4} = -\frac{\partial F_2}{\partial x_3} = -\frac{\partial F_3}{\partial x_2}. \quad (8)$$

*Proof.* The proof of this theorem can be analyzed in greater detail in [1].  $\square$

### 3. The Laplace's Equations

In this section we show that a new set of hypercomplex Laplace equations may be generated in four dimensions, through the use of Riemann-Cauchy like relations [1] Therefore, the functions that make up the quaternionic function, depend on  $x_1, x_2, x_3$  and  $x_4$  and are supposed to be of class  $C^2$  and thus the Schwartz's theorem is valid.

**Theorem 3.** *Let  $f(q)$  is an quaternionic function. If  $f(q)$  is of class  $C^2$  and satisfies the Cauchy-Riemann conditions, then*

$$\begin{aligned} \Delta f_1 &= 0, \\ \Delta f_2 &= 0, \\ \Delta f_3 &= 0, \\ \Delta f_4 &= 0. \end{aligned} \tag{9}$$

Demonstration: The first step to obtain the Laplace equation is the derivation of equations (5), (6), (7) and (8) over  $x_1, x_2, x_3$  and  $x_4$ . That will be done as follows: Firstly, in deriving the conditions of equation (5), we have:

$$\begin{aligned} \frac{\partial^2 F_1}{\partial x_1^2} &= \frac{\partial^2 F_2}{\partial x_1 \partial x_2} = \frac{\partial^2 F_3}{\partial x_1 \partial x_3} = \frac{\partial^2 F_4}{\partial x_1 \partial x_4}, \\ \frac{\partial^2 F_1}{\partial x_1 \partial x_2} &= \frac{\partial^2 F_2}{\partial x_2^2} = \frac{\partial^2 F_3}{\partial x_2 \partial x_3} = \frac{\partial^2 F_4}{\partial x_2 \partial x_4}, \\ \frac{\partial x_3 \partial x_1}{\partial^2 F_1} &= \frac{\partial x_3 \partial x_2}{\partial^2 F_2} = \frac{\partial x_3^2}{\partial^2 F_3} = \frac{\partial x_4 \partial x_3}{\partial^2 F_4}, \\ \frac{\partial x_1 \partial x_4}{\partial^2 F_1} &= \frac{\partial x_4 \partial x_2}{\partial^2 F_2} = \frac{\partial x_4 \partial x_3}{\partial^2 F_3} = \frac{\partial x_4^2}{\partial^2 F_4}. \end{aligned} \tag{10}$$

In deriving the conditions of equation (6), we have:

$$\begin{aligned} \frac{\partial^2 F_2}{\partial x_1^2} &= -\frac{\partial^2 F_1}{\partial x_1 \partial x_2} = -\frac{\partial^2 F_3}{\partial x_1 \partial x_4} = \frac{\partial^2 F_4}{\partial x_1 \partial x_3}, \\ \frac{\partial^2 F_2}{\partial x_1 \partial x_2} &= -\frac{\partial^1 F_1}{\partial x_2^2} = -\frac{\partial^2 F_3}{\partial x_2 \partial x_4} = \frac{\partial^2 F_4}{\partial x_3 \partial x_2}, \\ \frac{\partial x_3 \partial x_1}{\partial^2 F_2} &= -\frac{\partial x_3 \partial x_2}{\partial^2 F_1} = -\frac{\partial x_3 \partial x_4}{\partial^2 F_3} = \frac{\partial x_3^2}{\partial^2 F_4}, \\ \frac{\partial x_4 \partial x_1}{\partial^2 F_2} &= -\frac{\partial x_4 \partial x_2}{\partial^2 F_1} = -\frac{\partial x_4^2}{\partial^2 F_3} = \frac{\partial x_4 \partial x_3}{\partial^2 F_4}. \end{aligned} \tag{11}$$

In deriving the conditions of equation (7), we obtain that:

$$\begin{aligned}
 \frac{\partial^2 F_3}{\partial x_1^2} &= -\frac{\partial^2 F_1}{\partial x_1 \partial x_3} = -\frac{\partial^2 F_2}{\partial x_1 \partial x_4} = \frac{\partial^2 F_4}{\partial x_1 \partial x_2}, \\
 \frac{\partial^2 F_3}{\partial x_1 \partial x_2} &= -\frac{\partial^2 F_1}{\partial x_2 \partial x_3} = -\frac{\partial^2 F_2}{\partial x_2 \partial x_4} = \frac{\partial^2 F_4}{\partial x_2^2}, \\
 \frac{\partial^2 F_3}{\partial x_3 \partial x_1} &= -\frac{\partial^2 F_1}{\partial x_3^2} = -\frac{\partial^2 F_2}{\partial x_4 \partial x_3} = \frac{\partial^2 F_4}{\partial x_3 \partial x_2}, \\
 \frac{\partial^2 F_3}{\partial x_1 \partial x_4} &= -\frac{\partial^2 F_1}{\partial x_4 \partial x_3} = -\frac{\partial^2 F_2}{\partial x_4^2} = \frac{\partial^2 F_4}{\partial x_4 \partial x_2}.
 \end{aligned} \tag{12}$$

And finally in deriving the conditions of equation (8), we have:

$$\begin{aligned}
 \frac{\partial^2 F_4}{\partial x_1^2} &= \frac{\partial^2 F_1}{\partial x_1 \partial x_4} = -\frac{\partial^2 F_2}{\partial x_1 \partial x_3} = -\frac{\partial^2 F_3}{\partial x_1 \partial x_2}, \\
 \frac{\partial^2 F_4}{\partial x_1 \partial x_2} &= \frac{\partial^2 F_1}{\partial x_2 \partial x_4} = -\frac{\partial^2 F_2}{\partial x_2 \partial x_3} = -\frac{\partial^2 F_3}{\partial x_2^2}, \\
 \frac{\partial^2 F_4}{\partial x_3 \partial x_1} &= \frac{\partial^2 F_1}{\partial x_3 \partial x_4} = -\frac{\partial^2 F_2}{\partial x_3^2} = -\frac{\partial^2 F_3}{\partial x_3 \partial x_2}, \\
 \frac{\partial^2 F_4}{\partial x_1 \partial x_4} &= \frac{\partial^2 F_1}{\partial x_4^2} = -\frac{\partial^2 F_2}{\partial x_4 \partial x_3} = -\frac{\partial^2 F_3}{\partial x_4 \partial x_2}.
 \end{aligned} \tag{13}$$

In correlating those groups of partial derivatives (9), (10), (11) and (12), immediately follow the Laplace Equations:

$$\frac{\partial^2 F_1}{\partial x_1^2} + \frac{\partial^2 F_1}{\partial x_2^2} + \frac{\partial^2 F_1}{\partial x_3^2} + \frac{\partial^2 F_1}{\partial x_4^2} = 0, \tag{14}$$

$$\frac{\partial^2 F_2}{\partial x_1^2} + \frac{\partial^2 F_2}{\partial x_2^2} + \frac{\partial^2 F_2}{\partial x_3^2} + \frac{\partial^2 F_2}{\partial x_4^2} = 0, \tag{15}$$

$$\frac{\partial^2 F_3}{\partial x_1^2} + \frac{\partial^2 F_3}{\partial x_2^2} + \frac{\partial^2 F_3}{\partial x_3^2} + \frac{\partial^2 F_3}{\partial x_4^2} = 0, \tag{16}$$

and

$$\frac{\partial^2 F_4}{\partial x_1^2} + \frac{\partial^2 F_4}{\partial x_2^2} + \frac{\partial^2 F_4}{\partial x_3^2} + \frac{\partial^2 F_4}{\partial x_4^2} = 0. \tag{17}$$

Therefore, in a more simplified manner, the set of equations appears below:

$$\begin{aligned}
 \Delta f_1 &= 0, \\
 \Delta f_2 &= 0, \\
 \Delta f_3 &= 0, \\
 \Delta f_4 &= 0,
 \end{aligned} \tag{18}$$

where,

$$\begin{aligned}f_1 &= f_1(x_1, x_2, x_3, x_4), \\f_2 &= f_2(x_1, x_2, x_3, x_4), \\f_3 &= f_3(x_1, x_2, x_3, x_4), \\f_4 &= f_4(x_1, x_2, x_3, x_4).\end{aligned}\tag{19}$$

#### 4. Conclusion

In this note it is showed the feasibility of obtaining the equations of Laplace through the Cauchy-Riemann conditions for quaternions. This fact will allow the relationship between equations that can explain many physical phenomena. You can also use the above equations as a way of stating a theorem for harmonic functions for quaternions that satisfy the conditions of Cauchy. It is worth mentioning the importance of [1] in performing all relations used in this work.

#### References

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