

APPROXIMATION BY GENERALIZED FABER SERIES IN BERS SPACES ON QUASIDISKS

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Abstract: In this paper, by the conformal natural reflection introduced by Earle and Nag, we establish an integral representation for Bers space $B(D)$ where D is a quasidisk, and then we define generalized Faber series of functions in $B(D)$. We find a subspace $B_1(D)$ of $B(D)$ such that each $\varphi \in B_1(D)$ converges uniformly on compact subsets of D . Also, if a series $\sum_{m=3}^{\infty} a_m F_m'''(z)$ is convergent to $\varphi \in B_1(D)$ in the norm $\|\cdot\|_D$, we show that the a_m are the generalized Faber coefficients $a_m(\varphi)$ of φ .

AMS Subject Classification: 30C62, 30C10

Key Words: Faber series, Bers space, quasidisk, reproducing formula, conformally natural reflection

1. Introduction

Let D be a quasidisk in the extended complex plane $\overline{\mathbb{C}}$ i.e., D is the image of the open unit disk Δ under the quasiconformal self mapping of $\overline{\mathbb{C}}$. We denote by $B(D)$ the Bers space of holomorphic functions $\varphi(z)$ in D such that

$$\|\varphi\|_D = \sup\{|\varphi(z)|\lambda_D^{-2}(z) : z \in D\} < \infty,$$

where λ_D is the hyperbolic density on D (If $\infty \in D$ and $\varphi \in B(D)$, then $\varphi(z)$ has a zero of order ≥ 4 at ∞). It is well known that the space $B(D)$ is regarded as the tangent space of the Teichmüller space $T(D)$ (see [2], [7], [8], [9], [11], [12], [14], [15]). Due to the importance of the Bers space $B(D)$ in extremal quasiconformal mappings theory and Teichmüller spaces theory, it is meaningful to consider the polynomial approximation problem in $B(D)$.

Received: October 26, 2016

Revised: December 12, 2016

Published: March 19, 2017

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url: www.acadpubl.eu

Suppose $\Psi(w) \in \Sigma_0$ maps $\Delta^*(= \overline{C} \setminus \overline{\Delta})$ onto $D^*(= \overline{C} \setminus \overline{D})$, where Σ_0 is the class of univalent functions of the form $g(w) = w + b_1w^{-1} + b_2w^{-2} + \dots$ for $w \in \Delta^*$. It is known that \overline{D} is a compact subset of the closed ball $\{z : |z| \leq 2\}$ ([16]).

Faber polynomials play an important role in the study of univalent functions. Recall that the m th Faber polynomial F_m for $\Psi(w) \in \Sigma_0$ is a polynomial of degree m determined by the following expression:

$$\log \frac{\Psi(w) - z}{w} = - \sum_{m=1}^{\infty} \frac{1}{m} F_m(z) w^{-m}, z \in C, w \rightarrow \infty.$$

If $F_m(z)$ is the m th Faber polynomial for \overline{D} , then it is known that for every $(z, w) \in \overline{D} \times \Delta^*$,

$$\frac{\Psi'(w)}{\Psi(w) - z} = \sum_{m=0}^{\infty} \frac{F_m(z)}{w^{m+1}} \tag{1}$$

and the series converges uniformly and absolutely on compact subsets of $\overline{D} \times \Delta^*$. From (1), we get

$$\frac{6\Psi'(w)}{(\Psi(w) - z)^4} = \sum_{m=3}^{\infty} \frac{F_m'''(z)}{w^{m+1}}, (z, w) \in \overline{D} \times \Delta^* \tag{2}$$

where the series converges uniformly and absolutely on compact subsets of $\overline{D} \times \Delta^*$. More information for Faber polynomials and Faber expansions can be found in [10] and [17].

This paper is arranged as follows. In Section 2, by the conformal natural reflection introduced by Earle and Nag, we establish an integral representation for Bers space $B(D)$ where D is a quasidisk, and then we define a generalized Faber series of a function $\varphi \in B(D)$ to be the form $\sum_{m=3}^{\infty} a_m(\varphi) F_m'''(z)$. In Section 3, we define a subspace $B_1(D)$ of $B(D)$ such that each $\varphi \in B_1(D)$ converges uniformly on compact subsets of D . Also, if a series $\sum_{m=3}^{\infty} a_m F_m'''(z)$ is convergent to $\varphi \in B_1(D)$ in the norm $||| |||_D$, we show that the a_m are the generalized Faber coefficients $a_m(\varphi)$ of φ . Note that polynomial approximation problem has been studied by Cavus in Bergman spaces[4] and Israfilov in $A(\overline{D})$ where $A(\overline{D})$ denotes the class of functions which are analytic in D and continuous in \overline{D} [13] respectively.

2. Integral Representation of $\varphi \in B(D)$ by Conformal Natural Reflection

Theorem 1. [1] *If D is a Jordan domain, then D is a quasidisk if and only if there exists a quasiconformal reflection with respect to the boundary ∂D . When $\infty \in \partial D$, there exists a quasiconformal reflection satisfying a uniform Lipschitz condition.*

Theorem 2. [3] *Suppose D is a quasidisk with $\infty \in \partial D$. If h is a quasiconformal reflection with respect to ∂D and satisfies a uniform Lipschitz condition, then every $\varphi \in B(D)$ has the following reproducing formula*

$$\varphi(z) = -\frac{3}{\pi} \iint_{D^*} \frac{(\zeta - h(\zeta))^2}{(\zeta - z)^4} h_{\bar{\zeta}}(\zeta) \varphi(h(\zeta)) d\xi d\eta, z \in D. \tag{3}$$

In Theorem 1 Ahlfors gave a characterization of quasidisks by quasiconformal reflections and in Theorem 2 Bers established an integral representation of $\varphi \in B(D)$ where D is a quasidisk with $\infty \in \partial D$. We want to establish an integral representation of $\varphi \in B(D)$ where D is a quasidisk without the restriction $\infty \in \partial D$.

If $\infty \notin \partial D$, there exists a Möbius transformation $\gamma : \overline{C} \rightarrow \overline{C}$ such that $\infty \in \partial\gamma(D)$. By Theorem 1, there exists a quasiconformal reflection \hat{h} with respect to $\partial\gamma(D)$ and \hat{h} satisfies a uniform Lipschitz condition. Then $h = \gamma^{-1} \circ \hat{h} \circ \gamma$ is a quasiconformal reflection with respect to ∂D . By simple computations, we have the following

Theorem 3. *Suppose D is a quasidisk. There exists a quasiconformal reflection h with respect to ∂D such that every $\varphi \in B(D)$ has the following reproducing formula*

$$\varphi(z) = -\frac{3}{\pi} \iint_{D^*} \frac{(\zeta - h(\zeta))^2}{(\zeta - z)^4} h_{\bar{\zeta}}(\zeta) \varphi(h(\zeta)) d\xi d\eta, z \in D. \tag{4}$$

Earle and Nag introduced the concept of conformally natural quasiconformal reflection. Suppose S is a Jordan curve, D and D^* are two complementary components of S in \overline{C} . Two conformal mappings $f : \Delta \rightarrow D$ and $f^* : \Delta \rightarrow D^*$ determine the homeomorphism $f^{-1} \circ f^*$ and $f^{-1} \circ f^*$ has a conformally natural extension $ex(f^{-1} \circ f^*) : \overline{\Delta} \rightarrow \overline{\Delta}$ [5]. The mapping $j : \overline{C} \rightarrow \overline{C}$

$$j(z) = f^* \circ ex(f^{-1} \circ f^*)^{-1} \circ f^{-1}, z \in \overline{D}$$

$$j(z) = f \circ ex(f^{-1} \circ f^*) \circ (f^*)^{-1}, z \in \overline{D^*}$$

is called the conformally natural reflection of D . Earle and Nag[6] proved that j depends only on the domain D , independent of the choose of f and f^* ; j is a real analytic homeomorphism and is a quasiconformal reflection with respect to the boundary ∂D whenever D is a quasidisk; if j is the conformally natural reflection of D and $\gamma : \overline{C} \rightarrow \overline{C}$ is a Möbius transformation, then $\gamma \circ j \circ \gamma^{-1}$ is the conformally natural reflection of $\gamma(D)$. Earle, Gardiner and Lakic proved that for the quasidisk D with $\infty \in D$, the conformally natural quasiconformal reflection j satisfies a uniform Lipschitz condition. Hence by the proof of Theorem 3, we have the following result

Theorem 4. *Let D be a quasidisk and j be the conformally natural quasiconformal reflection with respect to the boundary ∂D . Then for every $\varphi \in B(D)$, we have*

$$\varphi(z) = -\frac{3}{\pi} \iint_{D^*} \frac{(\zeta - j(\zeta))^2}{(\zeta - z)^4} j_{\bar{\zeta}}(\zeta) \varphi(j(\zeta)) d\xi d\eta, z \in D. \tag{5}$$

Substituting $\zeta = \Psi(w)$ in (5), we get

$$\begin{aligned} \varphi(z) &= -\frac{3}{\pi} \iint_{\Delta^*} \frac{(\Psi(w) - j(\Psi(w)))^2}{(\Psi(w) - z)^4} j_{\bar{\zeta}}(\Psi(w)) \varphi(j(\Psi(w))) |\Psi'(w)|^2 dudv \\ &= -\frac{3}{\pi} \iint_{\Delta^*} \varphi(j(\Psi(w))) (\Psi(w) - j(\Psi(w)))^2 \overline{\Psi'(w)} j_{\bar{\zeta}}(\Psi(w)) \\ &\quad \cdot \frac{\Psi'(w)}{(\Psi(w) - z)^4} dudv. \end{aligned}$$

If we consider (2) and define coefficients $a_m(\varphi)$ by

$$a_m(\varphi) = -\frac{1}{2\pi} \iint_{\Delta^*} \frac{\varphi(j(\Psi(w))) (\Psi(w) - j(\Psi(w)))^2 \overline{\Psi'(w)}}{w^{m+1}} \cdot j_{\bar{\zeta}}(\Psi(w)) dudv, m = 3, 4, \dots \tag{6}$$

then we can associate a formal series $\sum_{m=3}^{\infty} a_m(\varphi) F_m'''(z)$ with $\varphi \in B(D)$, i.e.,

$$\varphi(z) \sim \sum_{m=3}^{\infty} a_m(\varphi) F_m'''(z). \tag{7}$$

We call this formal series a generalized Faber series of $\varphi \in B(D)$, and the coefficients $a_m(\varphi)$ are called generalized Faber coefficients of φ .

3. The Main Results and Proofs

Lemma 5. *The series $\sum_{m=3}^{\infty} \frac{|F_m'''(z)|^2}{m+1}$ is convergent uniformly on compact subsets of D .*

Proof. Let z be a fixed point in D . Then the power series $\sum_{m=3}^{\infty} \frac{F_m'''(z)}{m+1} w^{m+1}$ defines an analytic function $A(z, w)$ in Δ , i.e.,

$$A(z, w) := \sum_{m=3}^{\infty} \frac{F_m'''(z)}{m+1} w^{m+1}, w \in \Delta. \tag{8}$$

Thus, by taking derivative of (8) with respect to w and considering (2) we get

$$A'(z, w) = \sum_{m=3}^{\infty} F_m'''(z) w^m = \frac{6\Psi'(1/w)}{(\Psi(1/w) - z)^4 w}, w \in \Delta. \tag{9}$$

Let $0 < r < 1$. Since $\sum_{m=3}^{\infty} F_m'''(z) w^m$ is convergent uniformly and absolutely on the closed disc $\bar{B}(0, r)$, it follows that

$$\iint_{\bar{B}(0,r)} |A'(z, w)|^2 dudv = \pi \sum_{m=3}^{\infty} \frac{|F_m'''(z)|^2}{m+1} r^{2m+2}. \tag{10}$$

(9) and (10) show that

$$\pi \sum_{m=3}^{\infty} \frac{|F_m'''(z)|^2}{m+1} r^{2m+2} = 36 \iint_{\bar{B}(0,r)} \left| \frac{\Psi'(1/w)}{(\Psi(1/w) - z)^4 w} \right|^2 dudv. \tag{11}$$

Now we claim that

$$\iint_{\Delta^*} \left| \frac{\Psi'(w)}{(\Psi(w) - z)^4} \right|^2 dudv < +\infty.$$

Indeed by (2), we have

$$\lim_{w \rightarrow \infty} w^4 \frac{\Psi'(w)}{(\Psi(w) - z)^4} = 1.$$

Hence there exists some $R > 1$ such that $\left| \frac{\Psi'(w)}{(\Psi(w) - z)^4} \right| \leq \frac{2}{|w|^4}$ when $|w| > R$. Thus

$$\iint_{|w|>R} \left| \frac{\Psi'(w)}{(\Psi(w) - z)^4} \right|^2 dudv \leq \frac{4\pi}{3R^6}.$$

We choose some R' such that $|\Psi(w) - z| \leq R'$ whenever $z \in D$ and $1 < |w| < R$. Now for any $z \in D$, set $d(z, \partial D) = \inf\{|\zeta - z| : \zeta \in \partial D\}$ so that $|\Psi(w) - z| \geq d(z, \partial D)$ whenever $w \in \Delta^*$. Thus

$$\begin{aligned} \iint_{1<|w|<R} \left| \frac{\Psi'(w)}{(\Psi(w) - z)^4} \right|^2 dudv &= \iint_{\Psi(1<|w|<R)} \frac{1}{|\zeta - z|^8} d\xi d\eta \\ &\leq \iint_{d(z, \partial D) < |\zeta - z| < R'} \frac{1}{|\zeta - z|^8} d\xi d\eta \\ &\leq \frac{\pi}{3d^6(z, \partial D)}. \end{aligned}$$

Since

$$\lim_{r \rightarrow 1^-} \iint_{\bar{B}(0,r)} \left| \frac{\Psi'(1/w)}{(\Psi(1/w) - z)^4 w} \right|^2 dudv = \iint_{\Delta} \left| \frac{\Psi'(1/w)}{(\Psi(1/w) - z)^4 w} \right|^2 dudv$$

and

$$\iint_{\Delta} \left| \frac{\Psi'(1/w)}{(\Psi(1/w) - z)^4 w} \right|^2 dudv = \iint_{\Delta^*} \left| \frac{\Psi'(w)}{(\Psi(w) - z)^4} \right|^2 dudv < +\infty,$$

we get

$$\pi \sum_{m=3}^{\infty} \frac{|F_m'''(z)|^2}{m+1} = 36 \iint_{\Delta} \left| \frac{\Psi'(1/w)}{(\Psi(1/w) - z)^4 w} \right|^2 dudv.$$

On the other hand, it can be easily proved that

$$\iint_{\Delta} \left| \frac{\Psi'(1/w)}{(\Psi(1/w) - z)^4 w} \right|^2 dudv$$

is continuous in D , by Dini's theorem the series $\sum_{m=3}^{\infty} \frac{|F_m'''(z)|^2}{m+1}$ is convergent uniformly on compact subsets of D . The proof Lemma 5 is completed. \square

Definition 6. Denote by $B_1(D)$ the space of functions in $B(D)$ such that for every $\varphi \in B(D)$ there exists a constant M satisfying

$$\iint_{D^*} |\varphi(j(\zeta))|^2 |\zeta - j(\zeta)|^4 |j_{\bar{\zeta}}(\zeta)|^2 d\xi d\eta \leq M \|\varphi\|_D^2,$$

where j is the conformal natural quasiconformal reflection with respect to ∂D .

It is easy to see that $B_1(D)$ is a proper subspace of $B(D)$. Then we have the following

Theorem 7. Let $\varphi \in B_1(D)$. If $\sum_{m=3}^{\infty} a_m(\varphi) F_m'''(z)$ is a generalized Faber series of φ , then the series $\sum_{m=3}^{\infty} a_m(\varphi) F_m'''(z)$ converges uniformly to φ on compact subsets of D .

Proof. Let E be a compact subset of D . For every $z \in E$ and $\varphi \in B_1(D)$,

$$\begin{aligned} & \left| \varphi(z) - \sum_{m=3}^n a_m(\varphi) F_m'''(z) \right| \\ &= \frac{3}{\pi} \left| \iint_{\Delta^*} \varphi(j(\Psi(w))) (\Psi(w) - j(\Psi(w)))^2 \overline{\Psi'(w)} j_{\bar{\zeta}}(\Psi(w)) \left(\sum_{m=n+1}^{\infty} \frac{F_m'''(z)}{6w^{m+1}} \right) dudv \right| \\ &\leq \left(\iint_{\Delta^*} \left| \sum_{m=n+1}^{\infty} \frac{F_m'''(z)}{6w^{m+1}} \right|^2 dudv \right)^{\frac{1}{2}} \\ &\frac{3}{\pi} \left(\iint_{\Delta^*} |\varphi(j(\Psi(w)))|^2 |\Psi(w) - j(\Psi(w))|^4 |\Psi'(w)|^2 |j_{\bar{\zeta}}(\Psi(w))|^2 dudv \right)^{\frac{1}{2}} \\ &= \left(\iint_{\Delta^*} \left| \sum_{m=n+1}^{\infty} \frac{F_m'''(z)}{6w^{m+1}} \right|^2 dudv \right)^{\frac{1}{2}} \\ &\frac{3}{\pi} \left(\iint_{D^*} |\varphi(j(\zeta))|^2 |\zeta - j(\zeta)|^4 |j_{\bar{\zeta}}(\zeta)|^2 d\xi d\eta \right)^{\frac{1}{2}} \\ &\leq \frac{3\sqrt{M}\|\varphi\|_D}{\pi} \left(\iint_{\Delta^*} \left| \sum_{m=n+1}^{\infty} \frac{F_m'''(z)}{6w^{m+1}} \right|^2 dudv \right)^{\frac{1}{2}}. \end{aligned}$$

Let $1 < r < R < +\infty$.

$$\begin{aligned} \iint_{r < |w| < R} \left| \sum_{m=n+1}^{\infty} \frac{F_m'''(z)}{6w^{m+1}} \right|^2 dudv &= \frac{\pi}{36} \sum_{m=n+1}^{\infty} \frac{1}{m} \left(\frac{1}{r^{2m}} - \frac{1}{R^{2m}} \right) |F_m'''(z)|^2 \\ &\leq \frac{\pi}{9} \sum_{m=n+1}^{\infty} \frac{|F_m'''(z)|^2}{m+1}, \end{aligned}$$

and by letting $r \rightarrow 1^+$ and $R \rightarrow +\infty$ we get

$$\iint_{\Delta^*} \left| \sum_{m=n+1}^{\infty} \frac{F_m'''(z)}{6w^{m+1}} \right|^2 dudv \leq \frac{\pi}{9} \sum_{m=n+1}^{\infty} \frac{|F_m'''(z)|^2}{m+1}. \tag{12}$$

Therefore by Lemma 5, we conclude that $\sum_{m=3}^{\infty} a_m(\varphi) F_m'''(z)$ converges uniformly to φ on E . □

Corollary 8. Suppose $P_n(z) \in B_1(D)$ is a polynomial of degree n and $a_m(P_n)$ are its generalized Faber coefficients, then $a_m(P_n) = 0$ for all $m \geq n + 4$ and $P_n(z) = \sum_{m=3}^{n+3} a_m(P_n) F_m'''(z)$.

Proof. Let $z \in D$. By Theorem 7, we have

$$P_n(z) = \sum_{m=3}^{\infty} a_m(P_n) F_m'''(z).$$

It is obvious that $P_n(z)$ can be written in the form

$$P_n(z) = \sum_{v=3}^{n+3} A_v F_v'''(z).$$

Let j be the K -quasiconformal Earle-Nag reflection with respect to ∂D . Since j keeps ∂D pointwise fixed, $\Psi(w) = j(\Psi(w))$ with $|w| = 1$. It is easy to see that

$$\begin{aligned} & [F_v''(j(\Psi(w)))(\Psi(w) - j(\Psi(w)))^2]_{\bar{w}} \\ &= F_v'''(j(\Psi(w)))\overline{\Psi'(w)}j_{\bar{z}}(\Psi(w))(\Psi(w) - j(\Psi(w)))^2 \\ & - 2F_v''(j(\Psi(w)))\overline{\Psi'(w)}j_{\bar{z}}(\Psi(w))(\Psi(w) - j(\Psi(w))). \end{aligned}$$

and

$$\begin{aligned} & [F_v'(j(\Psi(w)))(\Psi(w) - j(\Psi(w)))]_{\bar{w}} \\ &= F_v''(j(\Psi(w)))\overline{\Psi'(w)}j_{\bar{z}}(\Psi(w))(\Psi(w) - j(\Psi(w))) \\ & - F_v'(j(\Psi(w)))\overline{\Psi'(w)}j_{\bar{z}}(\Psi(w)) \\ &= F_v''(j(\Psi(w)))\overline{\Psi'(w)}j_{\bar{z}}(\Psi(w))(\Psi(w) - j(\Psi(w))) \\ & - [F_v(j(\Psi(w)))]_{\bar{w}}. \end{aligned}$$

Hence we get by Green's formula

$$\begin{aligned} a_m(P_n) &= -\frac{1}{2\pi} \iint_{\Delta^*} \frac{P_n(j(\Psi(w)))(\Psi(w)-j(\Psi(w)))^2\overline{\Psi'(w)}}{w^{m+1}} j_{\bar{z}}(\Psi(w)) dudv \\ &= -\sum_{v=3}^{n+3} \frac{A_v}{2\pi} \iint_{\Delta^*} \frac{F_v'''(j(\Psi(w)))\overline{\Psi'(w)}(\Psi(w)-j(\Psi(w)))^2}{w^{m+1}} j_{\bar{z}}(\Psi(w)) dudv \\ &= -\sum_{v=3}^{n+3} \frac{A_v}{2\pi} \iint_{\Delta^*} \frac{[F_v''(j(\Psi(w)))(\Psi(w)-j(\Psi(w)))^2]_{\bar{w}}}{w^{m+1}} dudv \\ & - \sum_{v=3}^{n+3} \frac{A_v}{2\pi} \iint_{\Delta^*} \frac{2[F_v'(j(\Psi(w)))(\Psi(w)-j(\Psi(w)))]_{\bar{w}}}{w^{m+1}} dudv \\ & - \sum_{v=3}^{n+3} \frac{A_v}{2\pi} \iint_{\Delta^*} \frac{2[F_v(j(\Psi(w)))]_{\bar{w}}}{w^{m+1}} dudv \\ &= \sum_{v=3}^{n+3} \frac{A_v}{4\pi i} \int_{|w|=1} \frac{[F_v''(j(\Psi(w)))(\Psi(w)-j(\Psi(w)))^2]_{\bar{w}}}{w^{m+1}} dudv \\ & + \sum_{v=3}^{n+3} \frac{A_v}{4\pi i} \int_{|w|=1} \frac{2[F_v'(j(\Psi(w)))(\Psi(w)-j(\Psi(w)))]_{\bar{w}}}{w^{m+1}} dudv \\ & + \sum_{v=3}^{n+3} \frac{A_v}{4\pi i} \int_{|w|=1} \frac{2[F_v(j(\Psi(w)))]_{\bar{w}}}{w^{m+1}} dudv \\ &= \sum_{v=3}^{n+3} \frac{A_v}{2\pi i} \int_{|w|=1} \frac{[F_v(\Psi(w))]}{w^{m+1}} dudv. \end{aligned}$$

Since

$$\frac{1}{2\pi i} \int_{|w|=1} \frac{F_v(\Psi(w))}{w^{m+1}} dudv = \begin{cases} 1, & \text{if } v = m \\ 0, & \text{if } v \neq m, \end{cases} \tag{13}$$

see[10], it follows that $a_m(P_n) = A_m$, for $m = 3, \dots, n + 3$, and $a_m(P_n) = 0$, for $m \geq n + 4$. So $P_n(z) = \sum_{m=3}^{n+3} a_m(P_n) F_m'''(z)$. \square

Theorem 9. *Let $\{a_m\}$ be a complex number sequence. If the series $\sum_{m=3}^{\infty} a_m F_m'''(z)$ converges to a function $\varphi \in B_1(D)$ in the norm $\|\cdot\|_D$, then the a_m are the generalized Faber coefficients of φ .*

Proof. Let j be the K -quasiconformal Earle-Nag reflection with respect to ∂D , and $S_n(z) := \sum_{m=3}^{n+3} a_m F_m'''(z)$ be the n th partial sum of $\sum_{m=3}^{\infty} a_m F_m'''(z)$. Using (13), it can be shown that

$$\begin{aligned} & \lim_{n \rightarrow \infty} -\frac{1}{2\pi} \iint_{\Delta^*} \frac{S_n(j(\Psi(w)))(\Psi(w)-j(\Psi(w)))^2 \overline{\Psi'(w)}}{w^{m+1}} j_{\bar{\zeta}}(\Psi(w)) dudv \\ & = a_m, m = 3, 4, \dots, \end{aligned} \tag{14}$$

So, we get by using Hölder inequality

$$\begin{aligned} & |a_m(\varphi) - a_m| \\ & \leq \frac{1}{2\pi} \left| \iint_{\Delta^*} \frac{[\varphi(j(\Psi(w))) - S_n(j(\Psi(w)))](\Psi(w)-j(\Psi(w)))^2 \overline{\Psi'(w)}}{w^{m+1}} j_{\bar{\zeta}}(\Psi(w)) dudv \right| \\ & + \left| -\frac{1}{2\pi} \iint_{\Delta^*} \frac{S_n(j(\Psi(w)))(\Psi(w)-j(\Psi(w)))^2 \overline{\Psi'(w)}}{w^{m+1}} j_{\bar{\zeta}}(\Psi(w)) dudv - a_m \right| \\ & \leq \frac{1}{2\pi} \left(\iint_{\Delta^*} \frac{dudv}{|w|^{2m+2}} \right)^{\frac{1}{2}} \cdot \\ & \left(\iint_{\Delta^*} |\varphi(j(\Psi(w))) - S_n(j(\Psi(w)))|^2 |\Psi(w) \right. \\ & \left. - j(\Psi(w))\right|^4 |\Psi'(w)|^2 |j_{\bar{\zeta}}(\Psi(w))|^2 dudv \Big)^{\frac{1}{2}} \\ & + \left| -\frac{1}{2\pi} \iint_{\Delta^*} \frac{S_n(j(\Psi(w)))(\Psi(w)-j(\Psi(w)))^2 \overline{\Psi'(w)}}{w^{m+1}} j_{\bar{\zeta}}(\Psi(w)) dudv - a_m \right| \\ & \leq \sqrt{\frac{M}{12\pi}} \left(\iint_{D^*} |\varphi(j(\zeta)) - S_n(j(\zeta))|^2 |\zeta - j(\zeta)|^4 |j_{\bar{\zeta}}(\zeta)|^2 dudv \right)^{\frac{1}{2}} \\ & + \left| -\frac{1}{2\pi} \iint_{\Delta^*} \frac{S_n(j(\Psi(w)))(\Psi(w)-j(\Psi(w)))^2 \overline{\Psi'(w)}}{w^{m+1}} j_{\bar{\zeta}}(\Psi(w)) dudv - a_m \right| \\ & \leq \sqrt{\frac{M}{12\pi}} \|\varphi - S_n\|_D \\ & + \left| -\frac{1}{2\pi} \iint_{\Delta^*} \frac{S_n(j(\Psi(w)))(\Psi(w)-j(\Psi(w)))^2 \overline{\Psi'(w)}}{w^{m+1}} j_{\bar{\zeta}}(\Psi(w)) dudv - a_m \right|. \end{aligned}$$

$\lim_{n \rightarrow \infty} \|\varphi - S_n\|_D = 0$ and (14) show that $a_m(\varphi) = a_m$, and so the proof is completed. \square

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